

QCD Factorization for

Exclusive Non-leptonic B Decays

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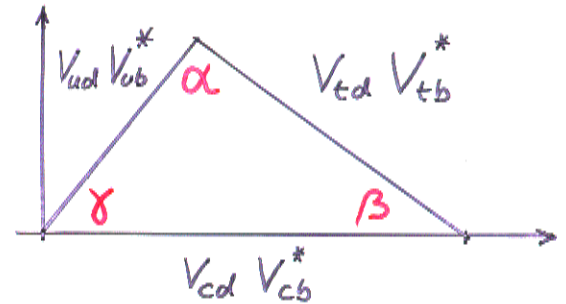
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Introduction

two-body nonleptonic B decays $B \rightarrow \pi\pi, \pi K, \dots$

- probe fundamental weak interactions of b-quarks
 \rightarrow CKM, CP violation, penguins, ...

$$A_{CP}(t) = \frac{\Gamma(B(t) \rightarrow \pi^+ \pi^-) - \Gamma(\bar{B}(t) \rightarrow \pi^+ \pi^-)}{\Gamma(B(t) \rightarrow \pi^+ \pi^-) + \Gamma(\bar{B}(t) \rightarrow \pi^+ \pi^-)}$$



pure direct CPV: $B^{\pm} \rightarrow K^{\pm} \pi^0, \dots$

- many channels $\pi, K, \rho, K^*, \phi, \dots$

- experimentally well accessible

CERN, Cornell, DESY,
FNAL, KEK, SLAC

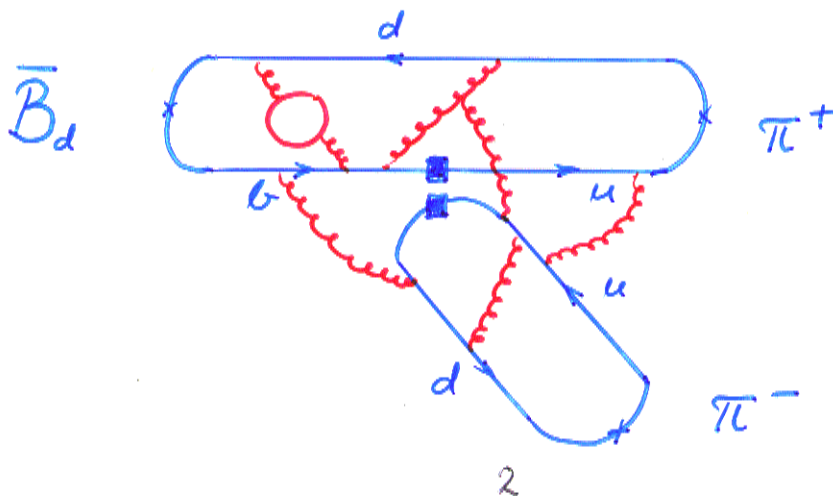
\Rightarrow rich phenomenology to study flavour physics

theory $\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \lambda_{\text{CKM}} C_i(\mu) Q_i$

problem: matrix elements $\langle \pi\pi | Q_i | B \rangle$

all hadronic final state

$\hookrightarrow (\bar{u}b)_{V-A} (d\bar{u})_{V-A}$



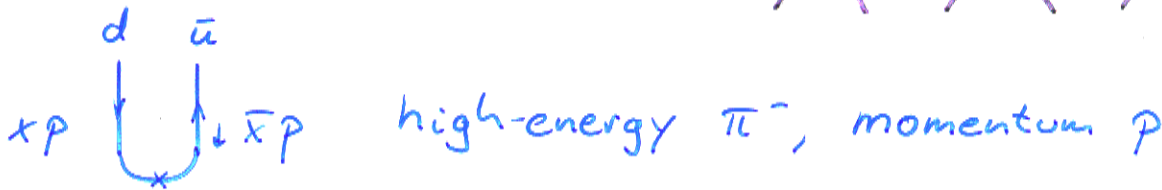
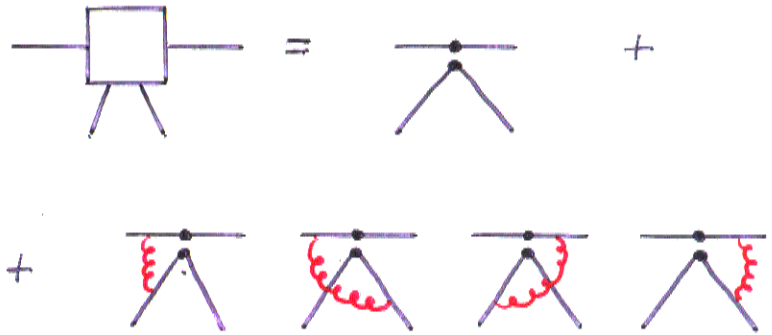
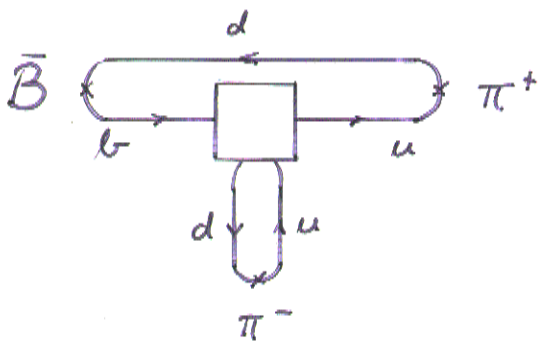
QCD factorization

$$m_B \gg \Lambda_{\text{QCD}}$$

$$A(B \rightarrow \pi\pi) \sim \underbrace{\langle \pi | j_1 | B \rangle \langle \pi | j_2 | 0 \rangle}_{\text{calculable in terms of universal quantities } f_+(0), f_\pi, \Phi_\pi, \dots} \left[1 + \mathcal{O}(\alpha_s) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{m_B}\right) \right]$$

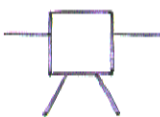
calculable in terms of universal quantities $f_+(0), f_\pi, \Phi_\pi, \dots$

compare: hard exclusive processes as



$$f_\pi \Phi_\pi(x) = f_\pi \cdot 6x\bar{x} \quad \bar{x} = 1-x$$

• $p_{\text{quark}} = x p$ to leading power

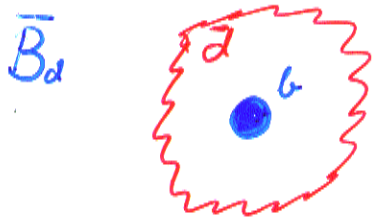
•  free of soft + collinear singularities
→ perturbative

• power suppression: end point region, higher Fock states

→ factorization formula

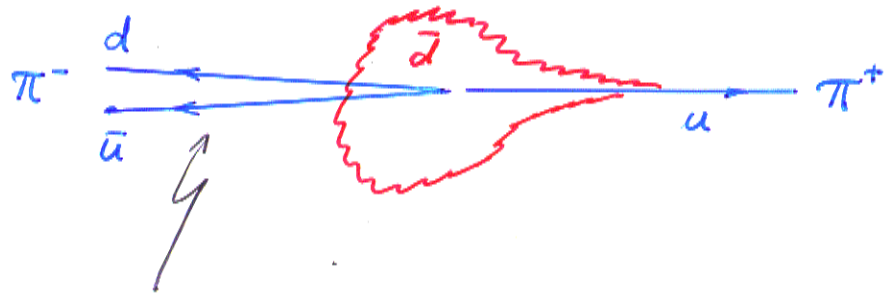
$$\langle \pi\pi | Q | B \rangle = F(B \rightarrow \pi) \cdot \int_0^1 dx T(x) \Phi_\pi(x)$$

intuitive picture

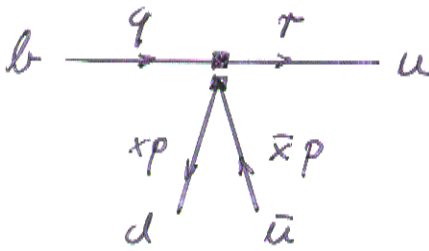


$$m_b \gg \Lambda_{QCD}$$

\Rightarrow



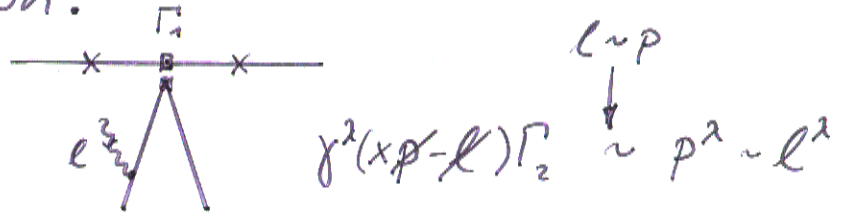
collinear, colour-singlet $\bar{u}d$ pair,
decoupled from IR interactions



soft cancellation:

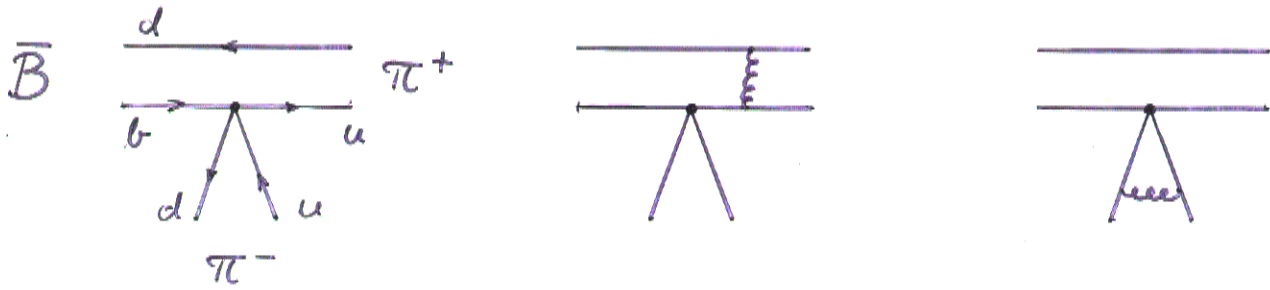
$$\begin{aligned}
 & \text{Diagram 1} + \text{Diagram 2} = \frac{\gamma^\lambda (x\not{p} + k) \Gamma_2}{(x\not{p} + k)^2} + \frac{\Gamma_2 (-\bar{x}\not{p} - k) \gamma^\lambda}{(\bar{x}\not{p} + k)^2} \\
 & \stackrel{k \text{ soft}}{=} \frac{2x \cdot p^\lambda \Gamma_2}{2x \cdot p \cdot k} + \frac{-2\bar{x} \cdot p^\lambda \Gamma_2}{2\bar{x} \cdot p \cdot k} = 0
 \end{aligned}$$

collinear cancellation:

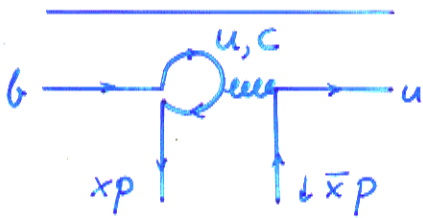


$$\begin{aligned}
 & \text{Diagram 1} + \text{Diagram 2} = \\
 & = \frac{\Gamma_1 (q - l + m_b) l}{(q - l)^2 - m_b^2} + \frac{l (r + l + m_u) \Gamma_1}{(r + l)^2 - m_u^2} = -\Gamma_1 + \Gamma_1 = 0
 \end{aligned}$$

$B \rightarrow \pi\pi$ - Overview



FSI \rightarrow phases

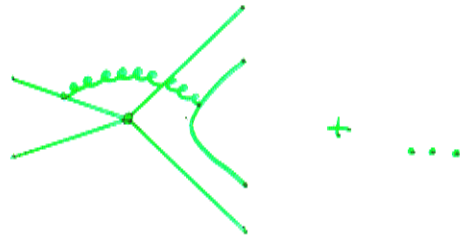


- gluon hard $k^2 = \bar{x} m_B^2$
- BSS mechanism
- $\langle G(k^2) \rangle = \int_0^1 dx G(\bar{x}) \Phi_\pi(x)$



gluon hard $k^2 \sim \Lambda m_B$

suppressed



$$\langle \pi\pi | Q | \bar{B} \rangle = F(B \rightarrow \pi) \cdot \int_0^1 dx T^I(x) \Phi_\pi(x) + \int_0^1 d\xi dx dy T^{II}(\xi, x, y) \Phi_B(\xi) \Phi_\pi(x) \Phi_\pi(y)$$

$$\xi = k_s^+ / p_B^+$$

$B \rightarrow \pi\pi$ amplitude from

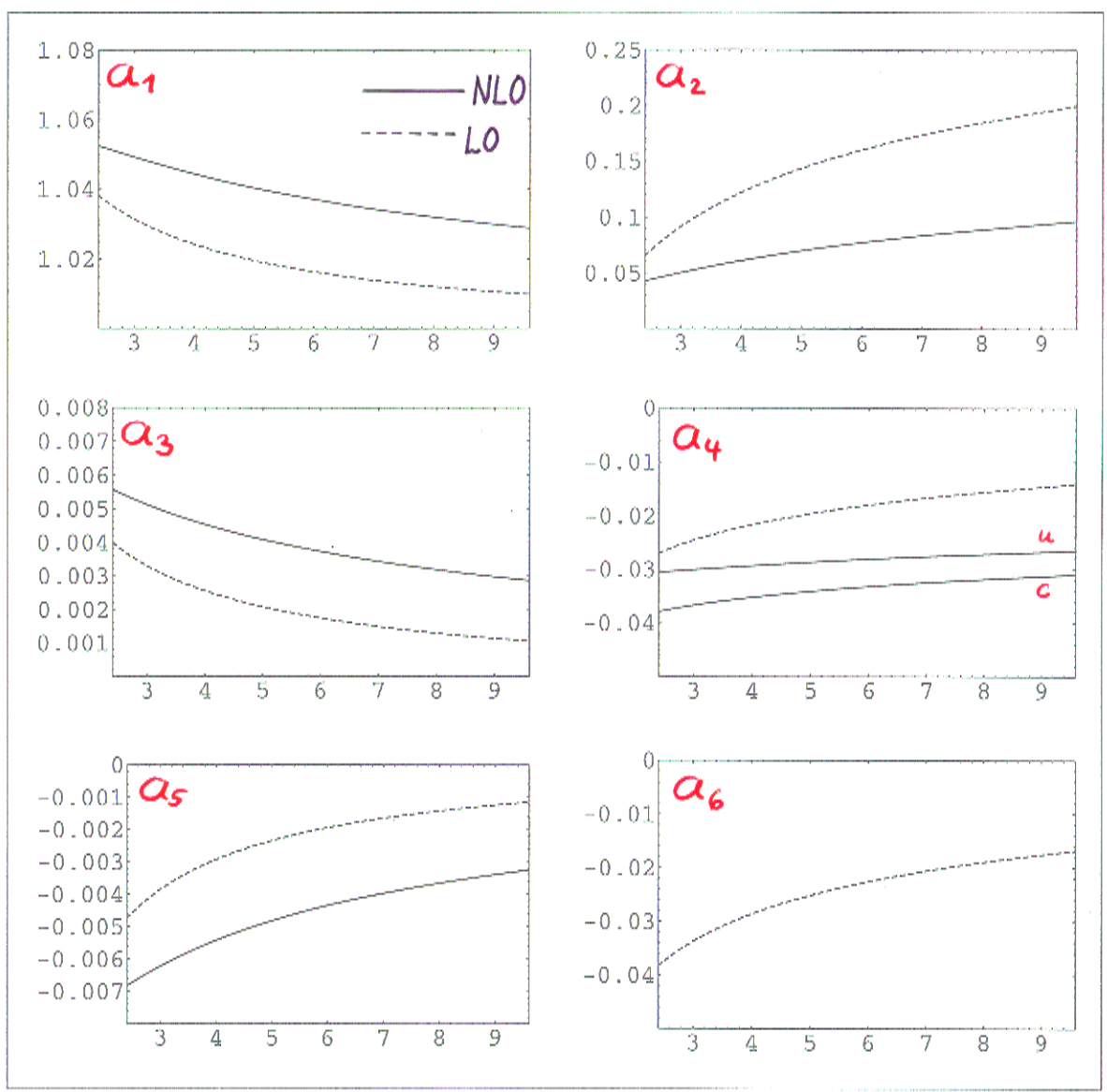
$$\hat{\mathcal{H}}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \sum_{p=u,c} \lambda_p \cdot$$

$$\begin{aligned} & \left[a_1^p (\bar{u}b)_{V-A} \otimes (\bar{d}u)_{V-A} + a_2^p (\bar{d}b)_{V-A} \otimes (\bar{u}u)_{V-A} + \right. \\ & + a_3 (\bar{d}b)_{V-A} \otimes (\bar{q}q)_{V-A} + a_4^p (\bar{q}b)_{V-A} \otimes (\bar{d}q)_{V-A} + \\ & \left. + a_5 (\bar{d}b)_{V-A} \otimes (\bar{q}q)_{V+A} + a_6^p (-2)(\bar{q}b)_{S-P} \otimes (\bar{d}q)_{S+P} \right] \end{aligned}$$

$$\begin{aligned} a_1^u = & C_1 + \frac{1}{3} C_2 + \frac{\alpha_s}{9\pi} \left(-12 \ln \frac{\mu}{m_b} - 18 + \int_0^1 g(x) \phi_2(x) dx \right) C_2 + \\ & + \frac{4\pi\alpha_s}{27} \frac{f_\pi f_B}{f_+ m_B^2} \int_0^1 \frac{\phi_\pi(x)}{x} \frac{\phi_\pi(y)}{y} dx dy \underbrace{\int_0^1 \frac{\phi_B(\xi)}{\xi} d\xi}_{\equiv m_B/\lambda_B} \cdot C_2 \end{aligned}$$

$$A(B \rightarrow \pi^+ \pi^-) \sim e^{i\gamma} + 0.23 e^{i\delta} \quad \delta = 6^\circ$$

$Re a_i$ ↑



→ μ/GeV

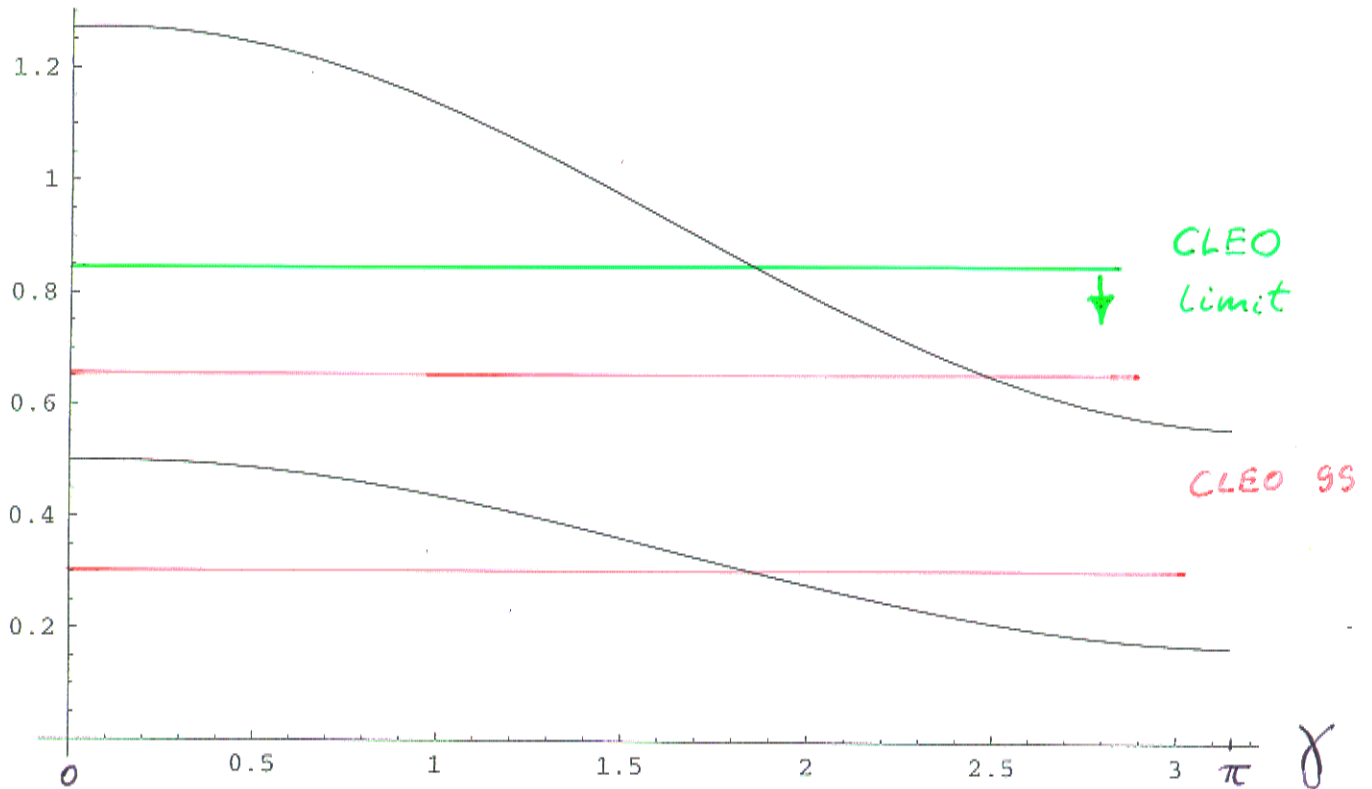
$$0.07 \leq \left| \frac{V_{ub}}{V_{cb}} \right| \leq 0.09$$

$$0.038 \leq |V_{cb}| \leq 0.040$$

$$0.25 \leq f_+(0) \leq 0.30 \quad \text{Bagan, Ball, Braun}$$

$$0.1 \leq \lambda_B / \text{GeV} \leq 0.4$$

$$B(B \rightarrow \pi^+ \pi^-) / 10^{-5}$$



$$B(B^+ \rightarrow \pi^+ \pi^0) = \left[\frac{f_+}{0.275} \frac{V_{cb}}{0.039} \frac{|V_{ub}/V_{cb}|}{0.08} \right]^2 (0.40 \pm 0.03) \cdot 10^{-5}$$

$$< 1.2 \cdot 10^{-5} \quad \text{CLEO}$$

$$B(B \rightarrow \pi^0 \pi^0) < 1.2 \cdot 10^{-6}$$

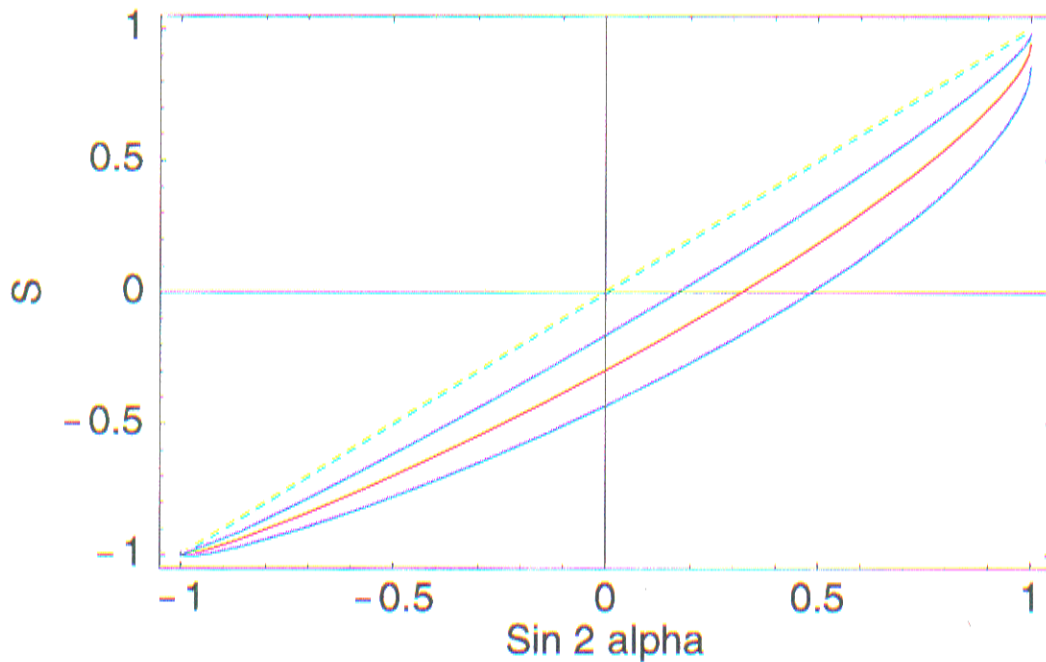
$\sin(2\alpha)$ from $B_d \rightarrow \pi^+ \pi^-$ decays

$$\mathcal{A}(t) = -S \cdot \sin(\Delta Mt) + C \cdot \cos(\Delta Mt)$$

$$S = \sin(2\alpha) + \mathcal{O}(P/T)$$

$$C = \mathcal{O}(P/T)$$

Fix $\sin(2\beta) = 0.7$.



green: no penguin amplitude (i.e. single weak phase contributes).

red: $a_6(\pi\pi)$ at leading order in α_s .

blue: twice $a_6(\pi\pi)$ and $a_6(\pi\pi) = 0$.

Summary

- factorization holds for $B \rightarrow \pi\pi$ in heavy quark limit

- calculated at LO and NLO in QCD

$$A(B \rightarrow \pi\pi) \sim f_{+}(0) f_{\pi} m_B^2 \left[1 + c \frac{1}{\ln \frac{m_B}{\Lambda}} + \mathcal{O}\left(\frac{\Lambda}{m_B}\right) \right]$$

\uparrow
non-factorizable, but calculable

- rescattering phases calculable

- other applications $B \rightarrow hh'$, $h, h' = \pi, \rho, K, K^*, \dots$

CPV $\bar{B} \rightarrow \pi^+ \pi^- \rightarrow \sin 2\alpha$

$B \rightarrow \pi K \rightarrow \gamma$

Fleischer, Mannel
Neubert, Rosner

$\bar{B} \rightarrow D^+ \pi^-$, $B^+ \rightarrow D^+ \pi^0$

Politzer, Wise

however: $\bar{B} \rightarrow D^0 \pi^0$ no factorization, but rel. power suppr.

- importance of power corrections ?

- QCD predictions for $B \rightarrow \pi\pi, \dots$ in HQ limit



to be tested in experiment