Physics 505 December, 2000 Ori Ganor

## Advanced topics in ordinary differential equations

## 13.1 $\Gamma$ -functions

The  $\Gamma$ -function is a generalization of the factorial to complex numbers.  $\Gamma(z)$  is defined for complex numbers z such that for positive integers:  $\Gamma(n) = (n-1)!$ . For  $\Re z > 0$  it is defined as:

$$\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt.$$

For  $\Re z \leq 0$  we have to use analytic continuation or the formula:

$$\Gamma(z) = \frac{1}{z}\Gamma(z+1).$$

This formula can be used several times so that:

$$\Gamma(z) = \frac{1}{z(z+1)\cdots(z+n-1)}\Gamma(z+n).$$

This formula shows that  $\Gamma(z)$  has simple poles for z=-n where n is a nonnegative integer. Note that:

$$\Gamma(\frac{1}{2}) = \int_0^\infty \frac{1}{\sqrt{t}} e^{-t} dt = \sqrt{\pi}.$$

## 13.2 Confluent hypergeometric functions

The confluent hypergeometric function is a holomorphic function defined by:

$$F(\alpha, \gamma, z) = 1 + \frac{\alpha}{\gamma} \frac{z}{1!} + \frac{\alpha(\alpha + 1)}{\gamma(\gamma + 1)} \frac{z^2}{2!} + \cdots$$

Note that it converges for all finte z. The parameter  $\alpha$  is arbitrary while  $\gamma$  is not allowed to be zero or a negative integer. Note also that  $F(\alpha, \alpha, z) = e^z$  and the relation:

$$F(\alpha, \gamma, z) = e^z F(\gamma - \alpha, \gamma, -z).$$

The function  $\psi(z) \equiv F(\alpha, \gamma, z)$  satisfies the differential equation:

$$z\psi'' + (\gamma - z)\psi' - \alpha\psi = 0.$$

The second independent solution of this equation is given by:  $z^{1-\gamma}F(\alpha-\gamma+1,2-\gamma,z)$ , as is easily checked. So the general solution is of the form:

$$\psi = C_1 F(\alpha, \gamma, z) + C_2 z^{1-\gamma} F(\alpha - \gamma + 1, 2 - \gamma, z).$$

We sometimes need to know the asymptotic expansion for large z. It is:

$$F(\alpha, \gamma, z) = \frac{\Gamma(\gamma)}{\Gamma(\gamma - \alpha)} (-z)^{-\alpha} G(\alpha, \alpha - \gamma + 1, -z) + \frac{\Gamma(\gamma)}{\Gamma(\alpha)} e^{z} z^{\alpha - \gamma} G(\gamma - \alpha, 1 - \alpha, z),$$

$$G(\alpha, \beta, z) = 1 + \frac{\alpha\beta}{1! z} + \frac{\alpha(\alpha + 1)\beta(\beta + 1)}{2! z^{2}} + \cdots$$

For  $\alpha$  that is a negative integer we have the equation:

$$F(-n,\gamma,z) = \frac{1}{\gamma(\gamma+1)\cdots(\gamma+n-1)} z^{1-\gamma} \frac{d^n}{dz^n} (e^{-z} z^{\gamma+n-1}).$$

## 13.3 Application to the Coulomb field

The confluent hypergeometric function appears in the solution of Schrödinger's equation in a Coulomb potential. The radial wave-function of the continuous spectrum is:

$$R_{kl}(r) = \frac{C_k}{(2l+1)!} (2kr)^l e^{-ikr} F(\frac{imZe^2}{\hbar^2 k} + l + 1, 2l + 2, 2ikr),$$

where  $C_k$  are normalization factors. If we take

$$C_k = \sqrt{\frac{2}{\pi}} k e^{\frac{\pi}{2k}} \left| \Gamma(l+1 - \frac{imZe^2}{\hbar^2 k}) \right|,$$

then the asymptotic behavior of  $R_{kl}$  can be calculated as:

$$R_{kl} pprox \sqrt{\frac{2}{\pi}} \frac{1}{r} \sin(kr + \frac{mZe^2}{\hbar^2 k} \log 2kr - \frac{1}{2}l\pi + \delta_l),$$

where:

$$\delta_l = \arg\Gamma(l+1 - \frac{imZe^2}{\hbar^2 k}).$$

Note that this is not quite the standard asymptotic behavior of a radial wave-function. The factor  $\log 2kr$  does not appear in the usual expansion. The Coulomb force does not decay fast enough as  $r \to \infty$  to have the usual asymptotic behavior. But since the only correction is a term that is independent of l, we can use  $\delta_l$  as the phase-shifts in our scattering formulas.

For more details see Landau& Lifshitz appendix d.