

New Lorentz Violating Nonlocal QFTs

Ori Ganor

UC Berkeley and LBNL

April 4, 2007

Talk at MIT

Based on

- OG, “A New Lorentz Violating Nonlocal Field Theory From String-Theory,” [[arXiv:hep-th/0609107](https://arxiv.org/abs/hep-th/0609107)]
- Aki Hashimoto, Sharon Jue, Bom Soo Kim, Anthony Ndirango and OG, “Aspects of Puff Field Theory,” [[arXiv:hep-th/0702030](https://arxiv.org/abs/hep-th/0702030)]

Introduction

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→ Nonlocality.

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- String theory construction?
- $c > 1$?

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- Bounds on photon dispersion relation:

$$\frac{\Delta c}{c} < \left| \frac{E}{10^{16} \text{GeV}} \right|, \quad E < 200 \text{KeV}.$$

[Amelino-Camelia, Ellis, Mavromatos, Nanopoulos and Sarkar, 1998]

$$\frac{\Delta c}{c} < \left| \frac{E}{10^{17} \text{GeV}} \right|, \quad 1 \text{MeV} < E < 17 \text{MeV}.$$

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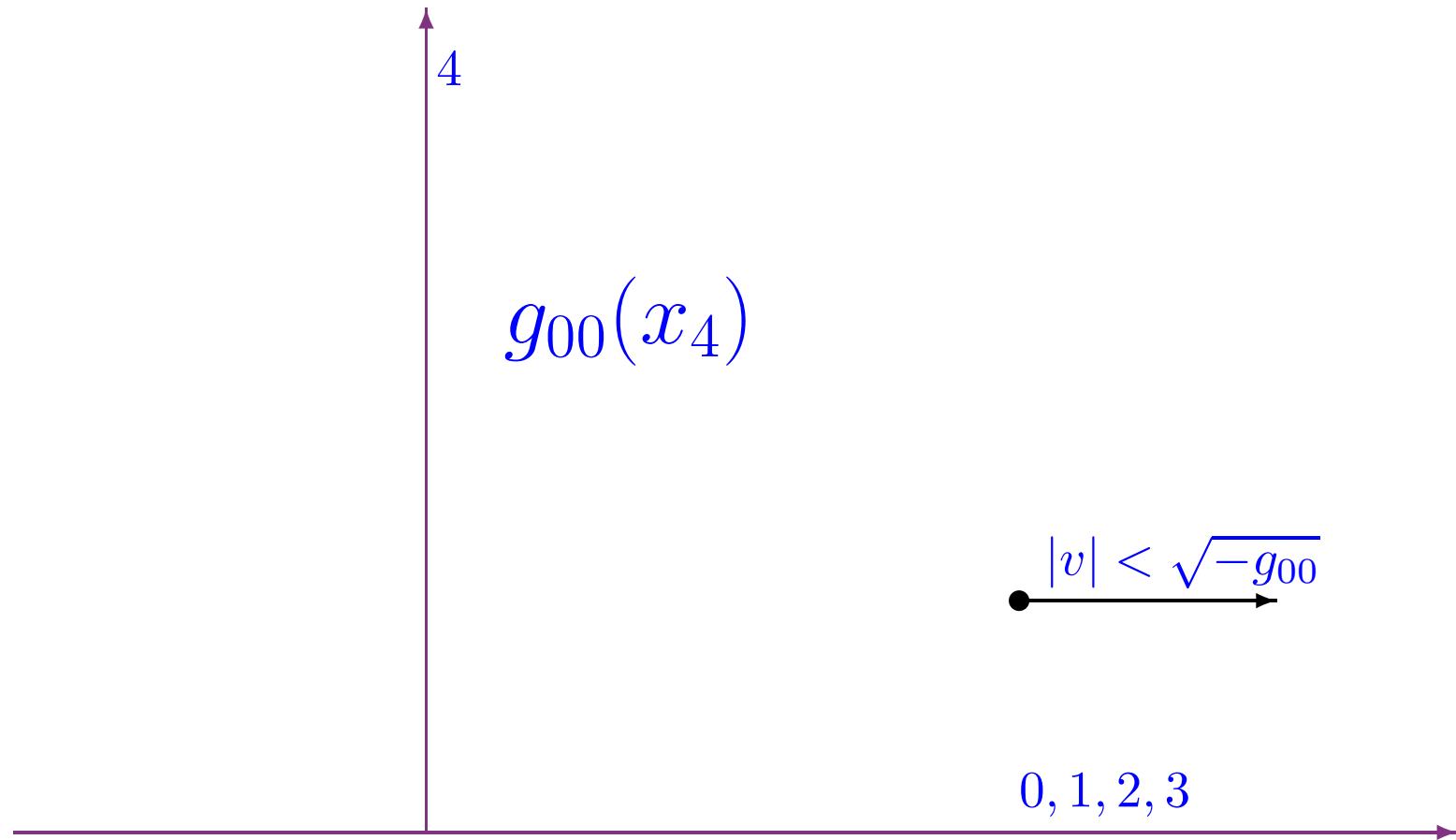
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Preserve Lorentz invariant photon dispersion relation.

Lorentz Violation & Extra Dimensions



[Csaki, Erlich and Grojean, 2000]

Nonlocality

- Yang-Mills theory on a Noncommutative \mathbb{R}^4 (NCYM)
[Douglas & Hull, Connes & Douglas & Schwarz, . . .]

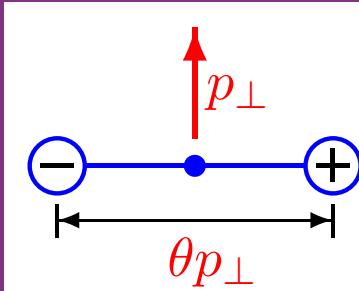
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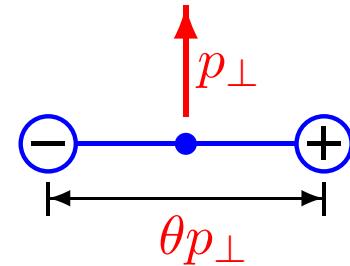
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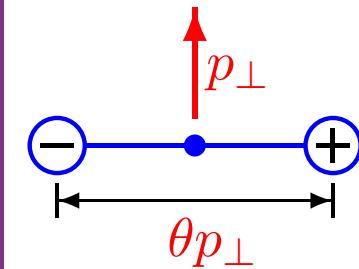
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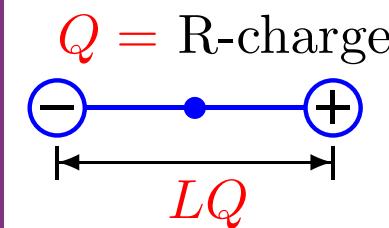
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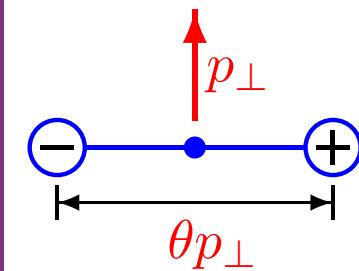
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- Dipole theories.



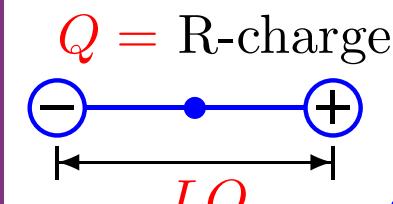
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$$D_\mu \Phi(x) =$$

$$\partial_\mu \Phi(x) - i A_\mu(x - \frac{1}{2}L) \Phi(x) + i \Phi(x) A_\mu(x + \frac{1}{2}L)$$

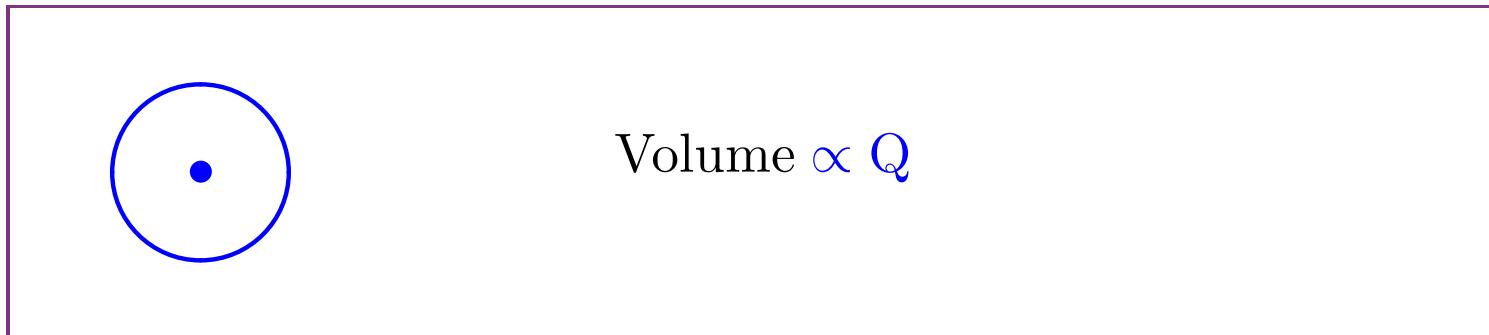
Nonlocality - Volume

- In this talk:



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Deformations of $N = 4$ SYM

$$L = L_{N=4} + \zeta \mathcal{O}^{(\Delta+4)} + \dots$$

$$[\zeta] = \text{Length}^\Delta$$

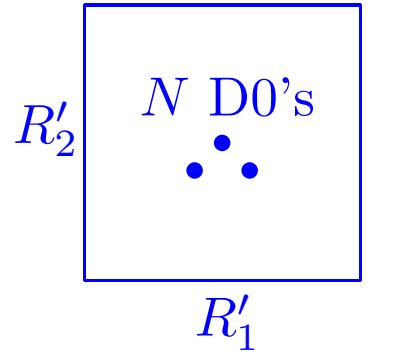
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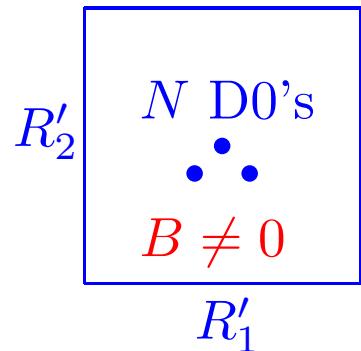
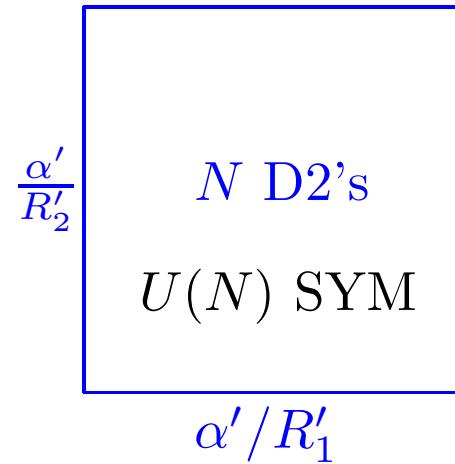
$$[\zeta] = \text{Length}^\Delta$$

Theory	Deformation $\zeta \mathcal{O}$	Δ	
NCYM	$(1/2g_{YM}^2)\theta^{\mu\nu} \text{Tr} \{ F_{\mu\sigma} F_{\nu\tau} F^{\sigma\tau} + F_{\sigma\tau} F^{\mu\nu} F_{\mu\nu} + \dots \}$	2	$\zeta \rightarrow \theta$
Dipole	$(1/2g_{YM}^2)L_{IJ}^\mu \text{Tr} \{ F_{\mu\nu} (\Phi^I D^\nu \Phi^J + \dots) + \dots \}$	1	$\zeta \rightarrow L$
This talk	3	

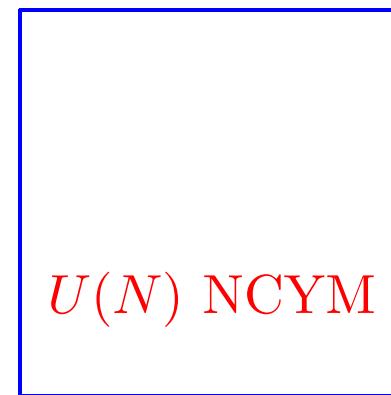
Douglas-Hull construction of NCYM



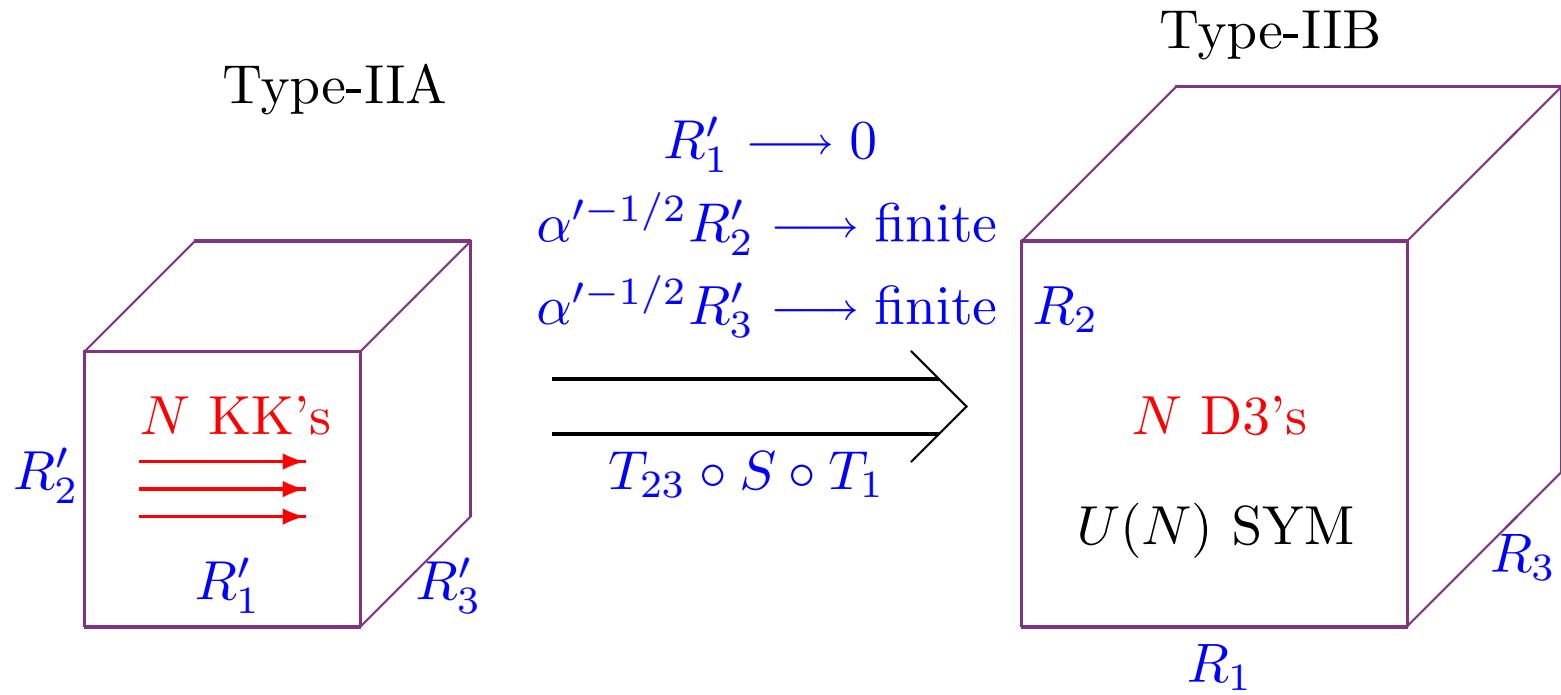
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Our Construction (1)



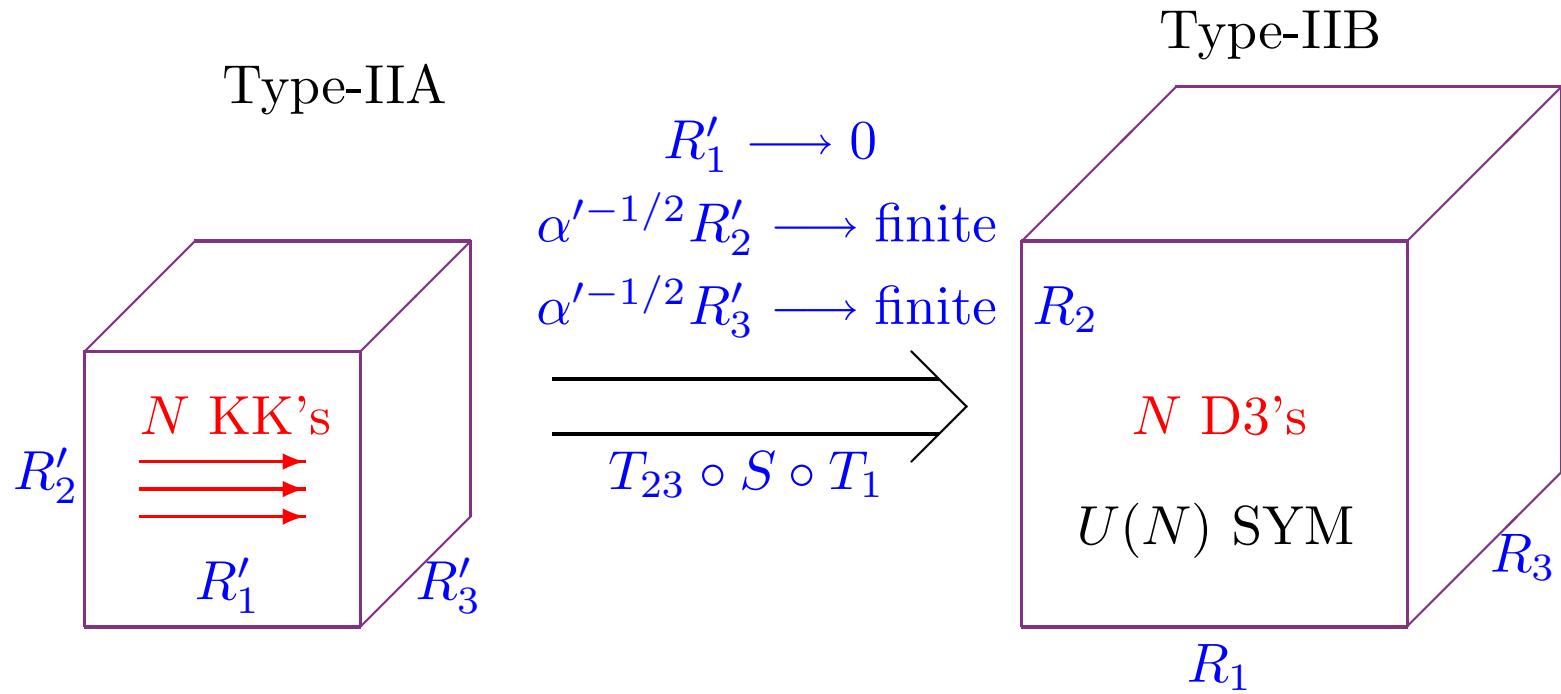
$$g_s = \alpha' {R'_3}^{-1} {R'_2}^{-1}$$

$$R_1 = \alpha'^{\frac{3}{4}} {g_s}'^{-\frac{1}{2}} {R'_1}^{-\frac{1}{2}}$$

$$R_2 = \alpha'^{\frac{5}{4}} {g_s}'^{\frac{1}{2}} {R'_1}^{-\frac{1}{2}} {R'_2}^{-1}$$

$$R_3 = \alpha'^{\frac{5}{4}} {g_s}'^{\frac{1}{2}} {R'_1}^{-\frac{1}{2}} {R'_3}^{-1}$$

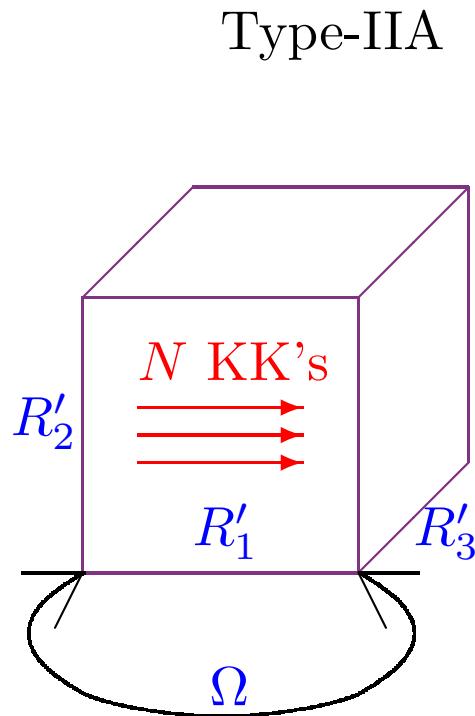
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$$g_s \rightarrow \text{finite} \quad \alpha'^{-1/2} R_k \rightarrow \infty, \quad (k = 1, 2, 3)$$

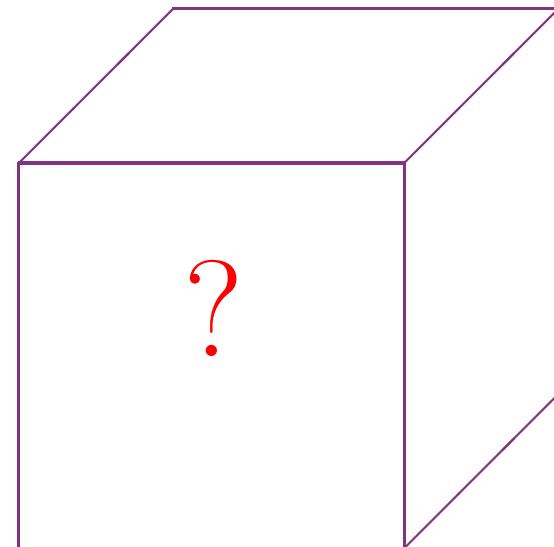
$$R_2/R_1, R_3/R_1 \rightarrow \text{finite}$$

Our Construction (2)



$$\begin{aligned} R'_1 &\longrightarrow 0 \\ \alpha'^{-1/2} R'_2 &\longrightarrow \text{finite} \\ \alpha'^{-1/2} R'_3 &\longrightarrow \text{finite} \end{aligned}$$

—————→



$$(x_1, x_2, x_3, \vec{y}) \sim (x_1 + 2\pi R_1, x_2, x_3, \Omega \vec{y})$$

$$\Omega \in \text{Spin}(6) \qquad \vec{y} \in \mathbb{R}^6$$

Scaling Ω

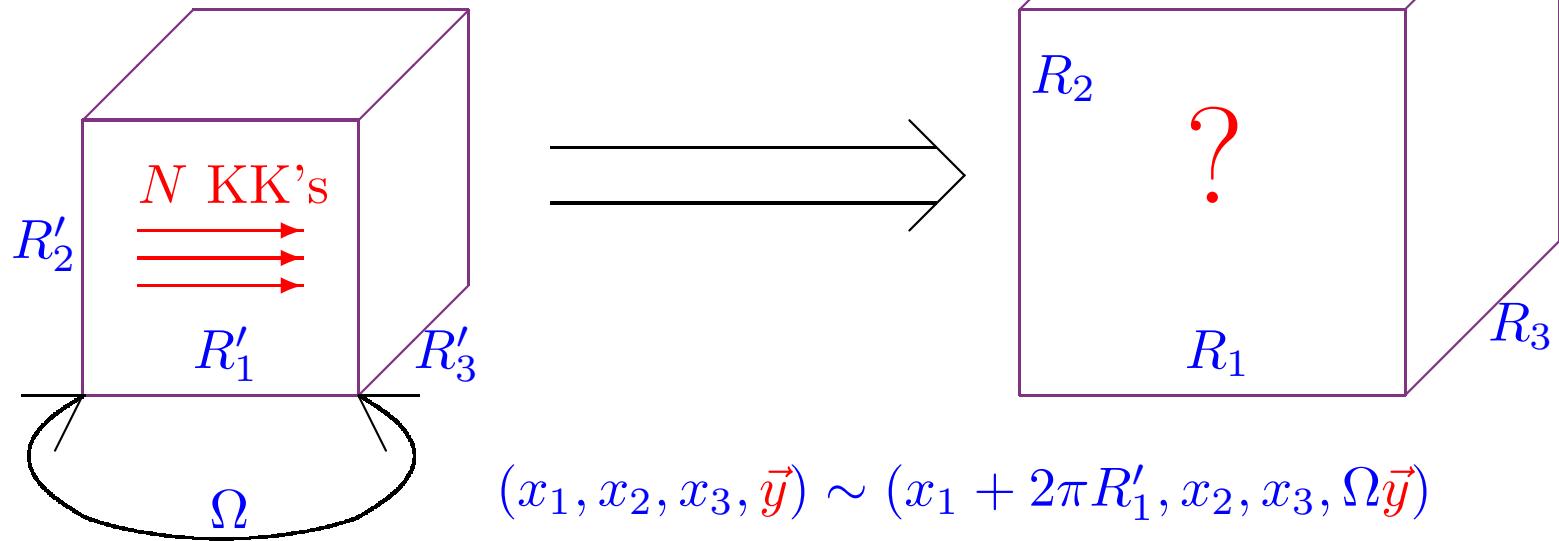
$$\Omega = \exp \left\{ \frac{2\pi}{R_1 R_2 R_3} \zeta \right\} \longrightarrow I, \quad \zeta \rightarrow \text{finite}$$

$$\zeta \equiv \begin{pmatrix} & & \beta_1 & & & \\ & -\beta_1 & & & & \\ & & & \beta_2 & & \\ & & -\beta_2 & & & \\ & & & & \beta_3 & \\ & & & & -\beta_3 & \end{pmatrix} \in so(6)$$

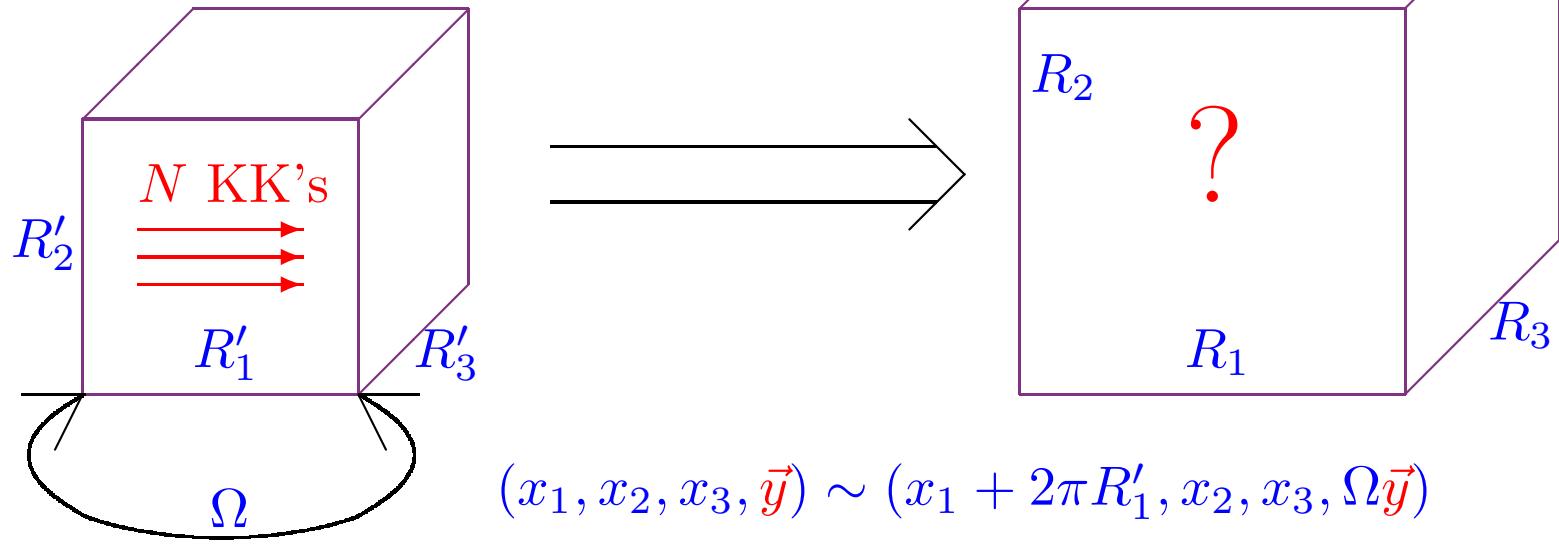
SUSY:

$$\begin{array}{ccccccccc} \{0\} & \subset & su(2) & \subset & su(3) & \subset & so(6) \\ \mathcal{N} = 4 & & \mathcal{N} = 2 & & \mathcal{N} = 1 & & \mathcal{N} = 0 \end{array}$$

Volume and R-charge

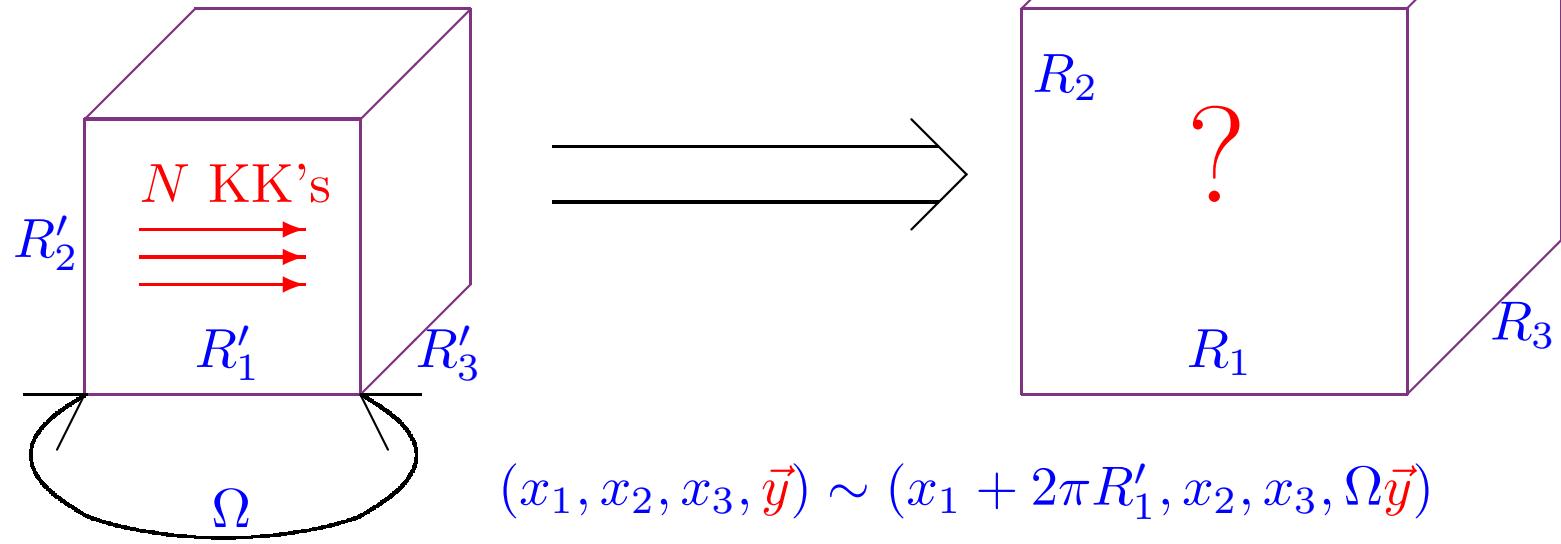


Volume and R-charge



$$z \equiv y_1 + iy_2, \quad \Omega = \exp \frac{4\pi i \beta}{R_1 R_2 R_3}, \quad f(x_1 + 2\pi R'_1, \Omega z) = f(x_1, z)$$

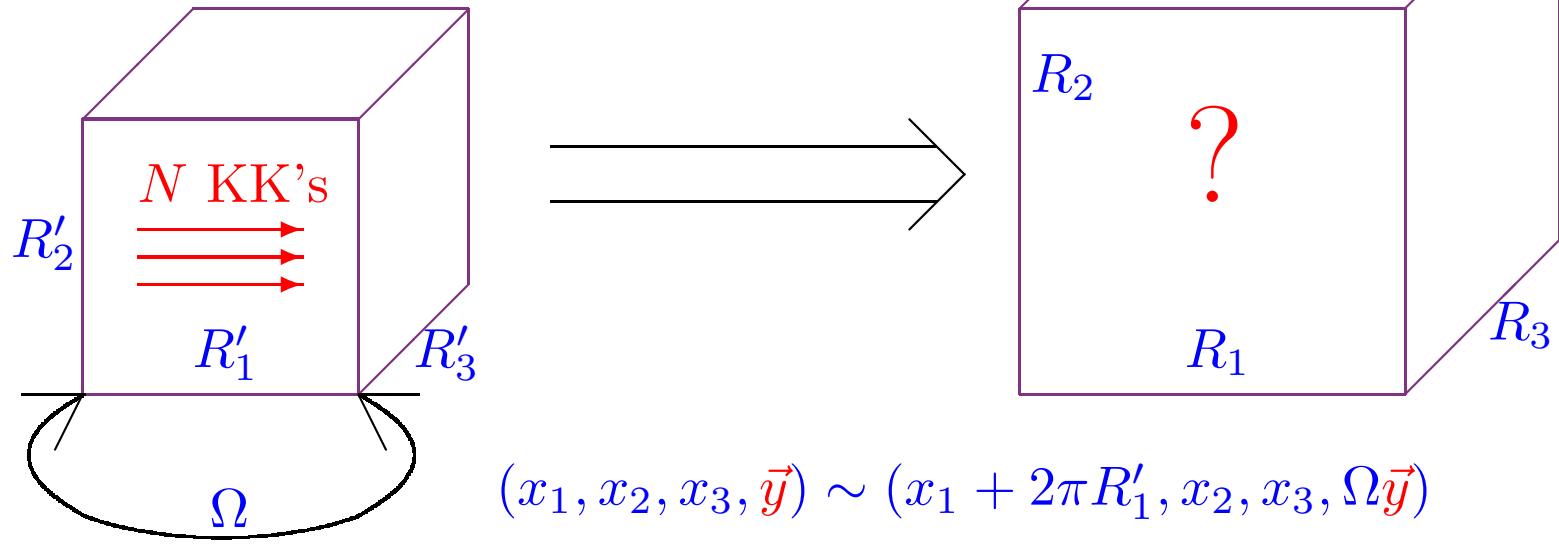
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$$f(x_1, z) = \sum_{p,j} C_{n,j}(|z|) z^{-j} e^{ipx_1/R'_1} \longrightarrow p \in \mathbb{Z} + \frac{2\beta j}{R_1 R_2 R_3}$$

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$$\Rightarrow \boxed{\text{“Vol(D3)”} = (2\pi)^3 (R_1 R_2 R_3 N + 2\beta j)}$$

Electric and Magnetic fluxes

$$V \equiv R_1 R_2 R_3,$$

$$\mathbf{P} \equiv \sum_{i=1}^3 \frac{k_i}{R_i} \hat{\mathbf{n}}_i, \quad \mathbf{E} \equiv \sum_{i=1}^3 \frac{e_i R_i}{2\pi V} \hat{\mathbf{n}}_i, \quad \mathbf{B} \equiv \sum_{i=1}^3 \frac{m_i R_i}{2\pi V} \hat{\mathbf{n}}_i.$$

$$\begin{aligned} E &= 2j\beta \frac{{M_s}^4}{g_s} + \frac{2\pi^2 V^2}{|NV+2j\beta|} \left(\frac{{g_{YM}}^2}{2\pi} \mathbf{E}^2 + \frac{2\pi}{{g_{YM}}^2} \mathbf{B}^2 \right) \\ &\quad + |\mathbf{P} - \frac{4\pi^2 V^2}{|NV+2j\beta|} \mathbf{E} \times \mathbf{B}|. \end{aligned}$$
$$\mathcal{N} = 2 \implies \beta \equiv \beta_1 = \beta_2, \quad \beta_3 = 0.$$

cf. NCSYM [Ho, Morariu & Zumino, Hofman & Verlinde, Konechny & Schwarz, Pioline & Schwarz, ... 1998-9].

Supergravity Dual

$$ds^2 = \frac{R^2}{r^2} K^{-\frac{1}{2}} \left[dx^2 + dy^2 + dz^2 - \left(dt - \frac{4\pi N}{r^2} \vec{n}^T \zeta d\vec{n} \right)^2 \right]$$

$$+ \frac{R^2}{r^2} K^{\frac{1}{2}} dr^2 + R^2 K^{\frac{1}{2}} d\Omega_5^2$$

$$\begin{aligned} C'_4 &= \frac{\pi N}{r^4} K^{-1} dt \wedge dx \wedge dy \wedge dz \\ &\quad - \frac{\pi N}{g_s \alpha'^2 r^6} K^{-1} \vec{n}^T \zeta d\vec{n} \wedge dx \wedge dy \wedge dz, \end{aligned}$$

$$K \equiv 1 + \frac{16\pi^2 N^2}{r^6} \vec{n}^T \zeta^T \zeta \vec{n}$$

$$\vec{n} \in S^5, \quad d\Omega_5^2 = \sum_{I=1}^6 dn_I^2, \quad R^4 \equiv 4\pi g_s N \alpha'^2$$

IR limit

$$\begin{aligned} ds^2 &= \frac{R^2}{r^2} [dr^2 + dx^2 + dy^2 + dz^2 - dt^2] + R^2 d\Omega_5^2 \\ &\quad + \frac{8\pi N R^2}{r^4} \vec{n}^T \zeta d\vec{n} dt + \dots \\ C'_4 &= \frac{\pi N}{r^4} dt \wedge dx \wedge dy \wedge dz \\ &\quad - \frac{\pi N}{g_s \alpha'^2 r^6} \vec{n}^T \zeta d\vec{n} \wedge dx \wedge dy \wedge dz + \dots \end{aligned}$$

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$$L_{N=4} \longrightarrow L_{N=4} + \zeta \mathcal{O} + \dots$$

Short multiplets

SYM Operator	desc	SUGRA	dim	spin	$SU(4)_R$
$\text{Tr } X^k$	-	$h_\alpha^\alpha \quad a_{\alpha\beta\gamma\delta}$	k	$(0, 0)$	$(0, k, 0)$
...
$\text{Tr } F_+ \bar{\lambda} \lambda X^k$	$Q^3 \bar{Q}$	$A_{\mu\alpha}$	$k + 5$	$(\frac{1}{2}, \frac{1}{2})$	$(1, k, 1)$
...
$\text{Tr } F_+ F_-^2 X^k$	$Q^4 \bar{Q}^2$	$A_{\mu\nu}$	$k + 6$	$(1, 0)$	$(0, k, 0)$
...
$\text{Tr } F_+ F_- \bar{\lambda} \lambda X^k$	$Q^3 \bar{Q}^3$	$h_{\mu\alpha} \quad a_{\mu\alpha\beta\gamma}$	$k + 7$	$(\frac{1}{2}, \frac{1}{2})$	$(1, k, 1)$
...
$\text{Tr } F_+^2 F_-^2 X^k$	$Q^4 \bar{Q}^4$	$h_\alpha^\alpha \quad a_{\alpha\beta\gamma\delta}$	$k + 8$	$(0, 0)$	$(0, k, 0)$

taken from [D'Hoker & Freedman, 2003]

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...
$\text{Tr } F_+^2 F_-^2 X^k$	$Q^4 \bar{Q}^4$	$h_\alpha^\alpha \quad a_{\alpha\beta\gamma\delta}$	$k + 8$	$(0, 0)$	$(0, k, 0)$

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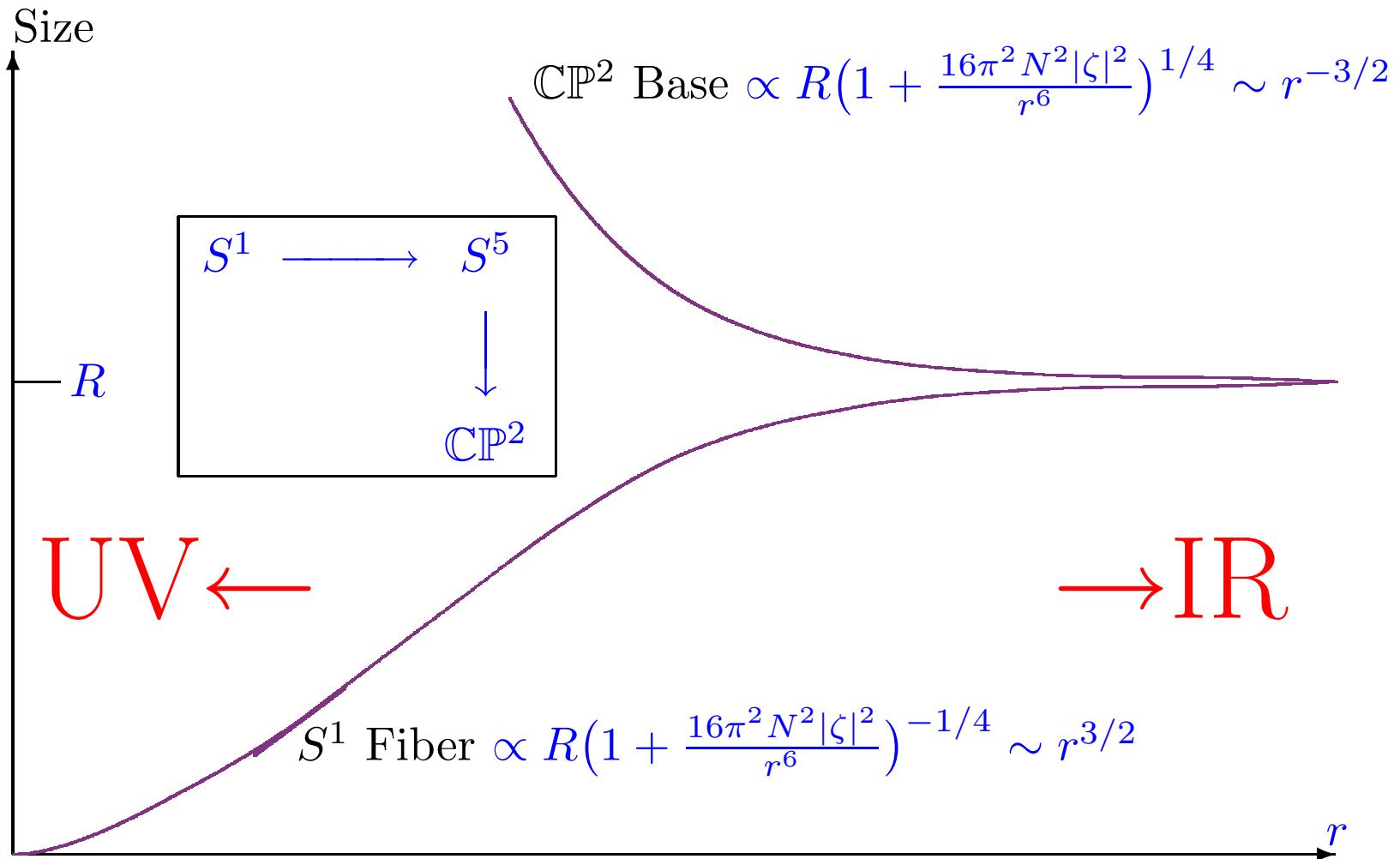
$$\zeta \mathcal{O} \rightarrow \zeta_{AB}^\mu T_{\mu\nu} J^{\nu AB} + \zeta_{AB}^\mu \epsilon^{ABCDEF} \epsilon_{\mu\nu\sigma\tau} X^C \partial_\nu X^D \partial_\sigma X^E \partial_\tau X^F + \dots$$

$J^{\nu AB}$ \equiv R-current

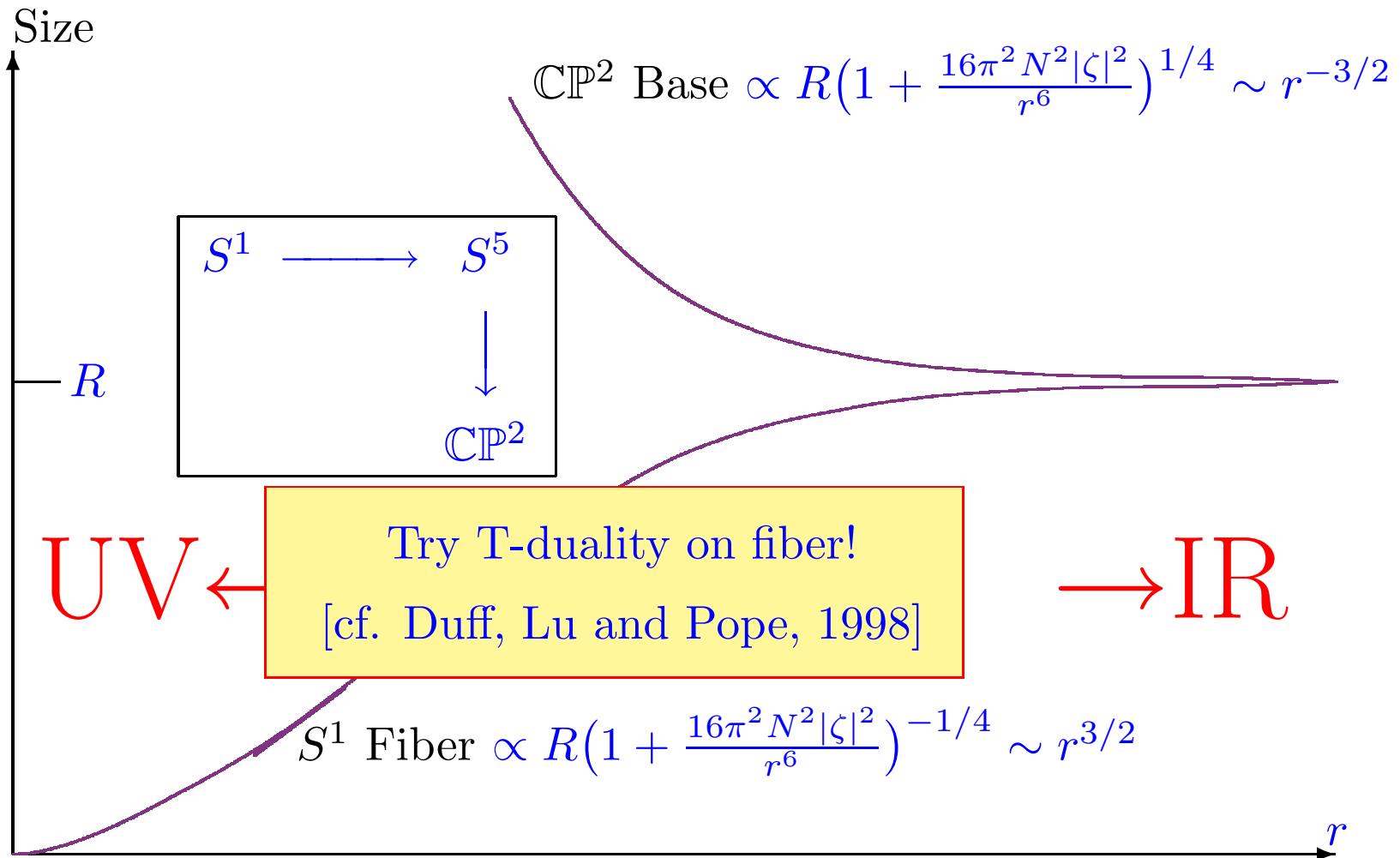
UV limit (1)

$$\begin{array}{ccc} S^1 & \longrightarrow & S^5 \\ & & \downarrow \\ & & \mathbb{CP}^2 \end{array}$$

UV limit (1)



UV limit (1)



UV Limit (2)

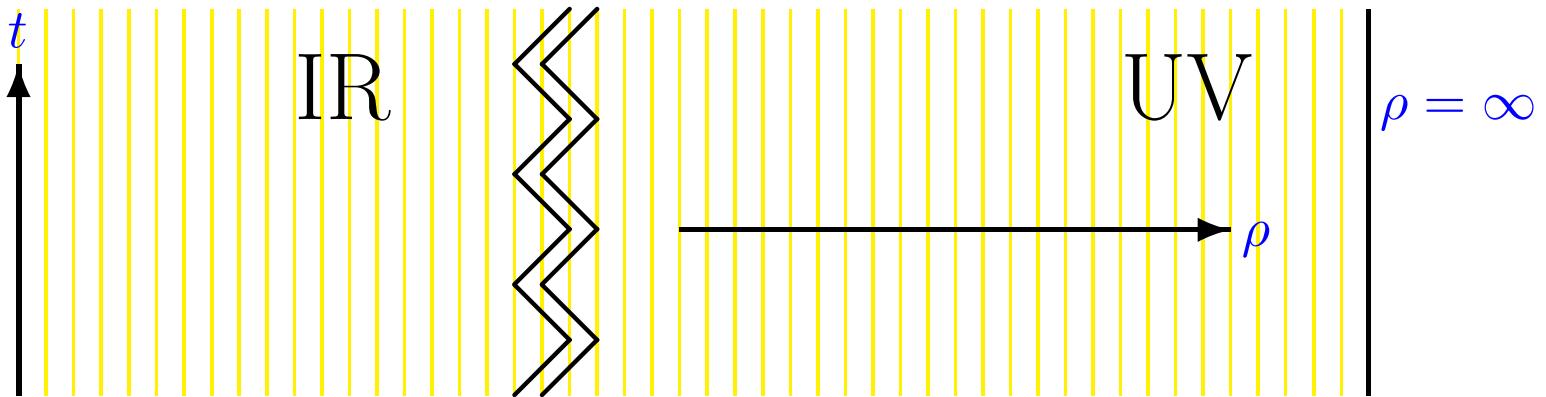
T-duality on fiber: Type-IIB \Rightarrow M.

$$\begin{aligned} \ell_P^{-2} ds_M^2 &= (4\pi N)^{-1/3} g_s^{-1} \rho^2 [\Delta^{-2/3} (dx^2 + dy^2 + dz^2) - \Delta^{1/3} dt^2] \\ &\quad + (4\pi N)^{2/3} \Delta^{1/3} \rho^{-2} d\rho^2 + (4\pi N)^{2/3} \Delta^{1/3} ds_B^2 \\ &\quad + (4\pi N)^{-1/3} \Delta^{1/3} (g_s^{-1} d\xi^2 + g_s d\eta^2), \end{aligned}$$

$$\begin{aligned} \ell_P^{-3} G_4 &= 2\pi N \omega \wedge \omega + 2\omega \wedge d\eta \wedge d\xi \\ &\quad + \frac{2}{3} (4\pi N)^{-2} g_s^{-3} \beta d\left(\frac{\rho^6}{\Delta}\right) \wedge dx \wedge dy \wedge dz, \end{aligned}$$

$$\Delta \equiv 1 + (4\pi N)^{-1} g_s^{-3} \beta^2 \rho^6, \quad (\rho \propto 1/r)$$

Extreme UV Limit



$$\begin{aligned} \ell_P^{-2} ds_M^2 &\approx -(4\pi N)^{-2/3} g_s^{-2} \beta^{2/3} \rho^4 dt^2 \\ &+ (4\pi N)^{1/3} g_s \beta^{-4/3} \rho^{-2} (dx^2 + dy^2 + dz^2) \\ &+ (4\pi N)^{1/3} g_s^{-1} \beta^{2/3} (d\rho^2 + \rho^2 ds_B^2) + (4\pi N)^{-2/3} \beta^{2/3} \rho^2 (g_s^{-2} d\xi^2 + d\eta^2), \end{aligned}$$

$$\ell_P^{-3} G_4 \approx 2\pi N \omega \wedge \omega + 2\omega \wedge d\eta \wedge d\xi - 4g_s^3 \beta^{-3} \rho^{-7} d\rho \wedge dx \wedge dy \wedge dz.$$

Extreme UV Limit (2)

$$\rho \rightarrow \infty$$

- Curvature $\rightarrow 0.$
- Geodesically complete.
- Redshift in frequency ($g_{00} \rightarrow \infty$)
- Blueshift in wavelength ($g_{xx}, g_{yy}, g_{zz} \rightarrow 0$)

Degrees of Freedom

UV regime of PFT – holographically dual to weakly-coupled background (?)

Is the spectrum discrete?

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For particle trajectories with fixed (E, \vec{p}) :

$$\rho \leq \rho_{\max} \equiv (4\pi N)^{1/6} g_s^{1/2} \beta^{-1/3} \left(\frac{E}{|\vec{p}|} \right)^{1/3}$$

Suggests a discrete spectrum!

cf. Little-String-Theory. [Aharony & Berkooz & Kutasov & Seiberg, 1998]

Effects of TJ coupling

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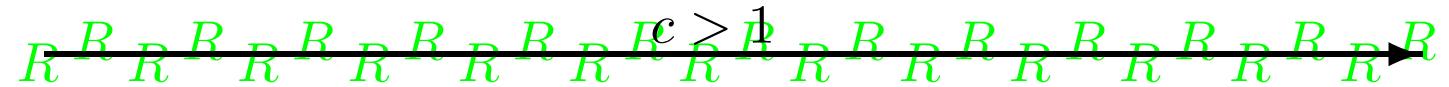
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$$g_{00} \rightarrow g_{00} - \zeta \langle J^0 \rangle + \dots, \quad v_{\max} \approx 1 + \frac{1}{2} \zeta \langle J^0 \rangle.$$



- In dense matter:

$$\Delta V_{\text{eff}} \approx \zeta \langle T^{00} \rangle Q^{(R)}, \quad Q^{(R)} \equiv \int J^0 d^3x = (\text{R-charge}).$$

Effects of TJ coupling

$$L = L_{(N=4)} + \zeta(T^{00}J_0 + \dots) + \dots$$

- In R-charged matter:

$$g_{00} \rightarrow g_{00} - \zeta \langle J^0 \rangle + \dots, \quad v_{\max} \approx 1 + \frac{1}{2} \zeta \langle J^0 \rangle.$$



- In dense matter:

$$\Delta V_{\text{eff}} \approx \zeta \langle T^{00} \rangle Q^{(R)}, \quad Q^{(R)} \equiv \int J^0 d^3x = (\text{R-charge}).$$

Trap R-charged particles?

More on $c > 1$

Extreme UV:

$$\begin{aligned}\ell_P^{-2} ds_M^2 \approx & -(4\pi N)^{-2/3} g_s^{-2} \beta^{2/3} \rho^4 dt^2 \\ & +(4\pi N)^{1/3} g_s \beta^{-4/3} \rho^{-2} (dx^2 + dy^2 + dz^2) \\ & +(4\pi N)^{1/3} g_s^{-1} \beta^{2/3} d\rho^2 + \dots\end{aligned}$$

More on $c > 1$

Extreme UV:

$$\begin{aligned}\ell_P^{-2} ds_M^2 \approx & -(4\pi N)^{-2/3} g_s^{-2} \beta^{2/3} \rho^4 dt^2 \\ & +(4\pi N)^{1/3} g_s \beta^{-4/3} \rho^{-2} (dx^2 + dy^2 + dz^2) \\ & +(4\pi N)^{1/3} g_s^{-1} \beta^{2/3} d\rho^2 + \dots\end{aligned}$$

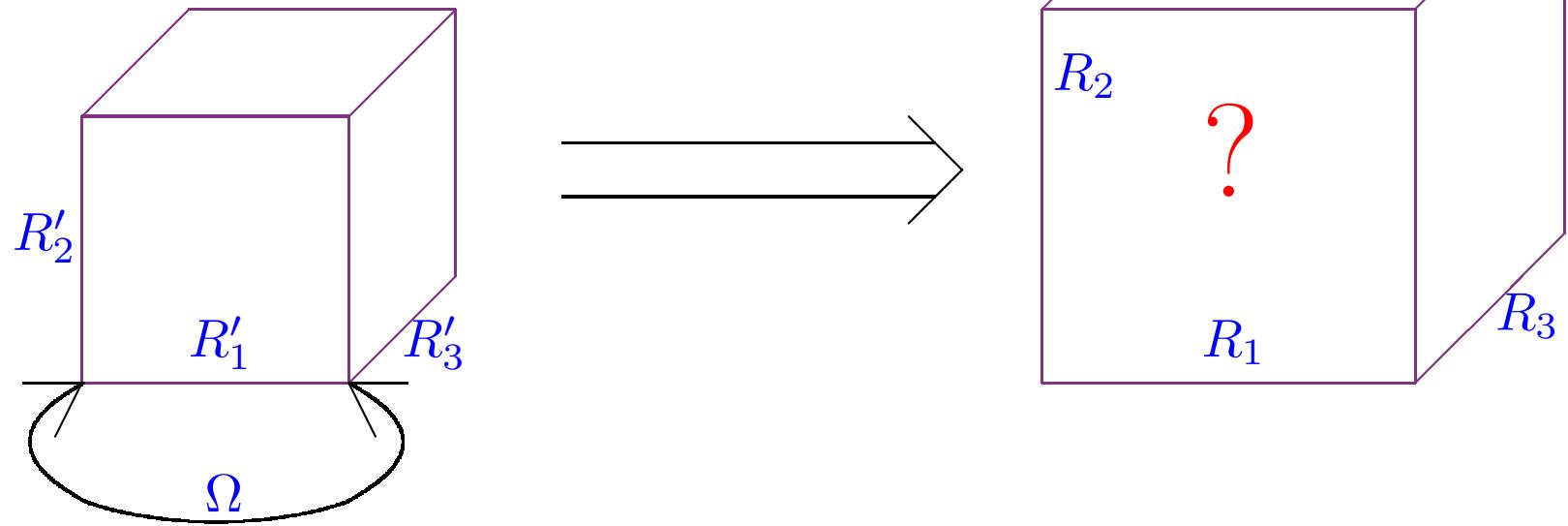
For a particle of mass m :

$$-\ell_P^{-2} m^2 \geq -(4\pi N)^{2/3} g_s^{-2} \beta^{-2/3} \frac{E^2}{\rho^4}.$$

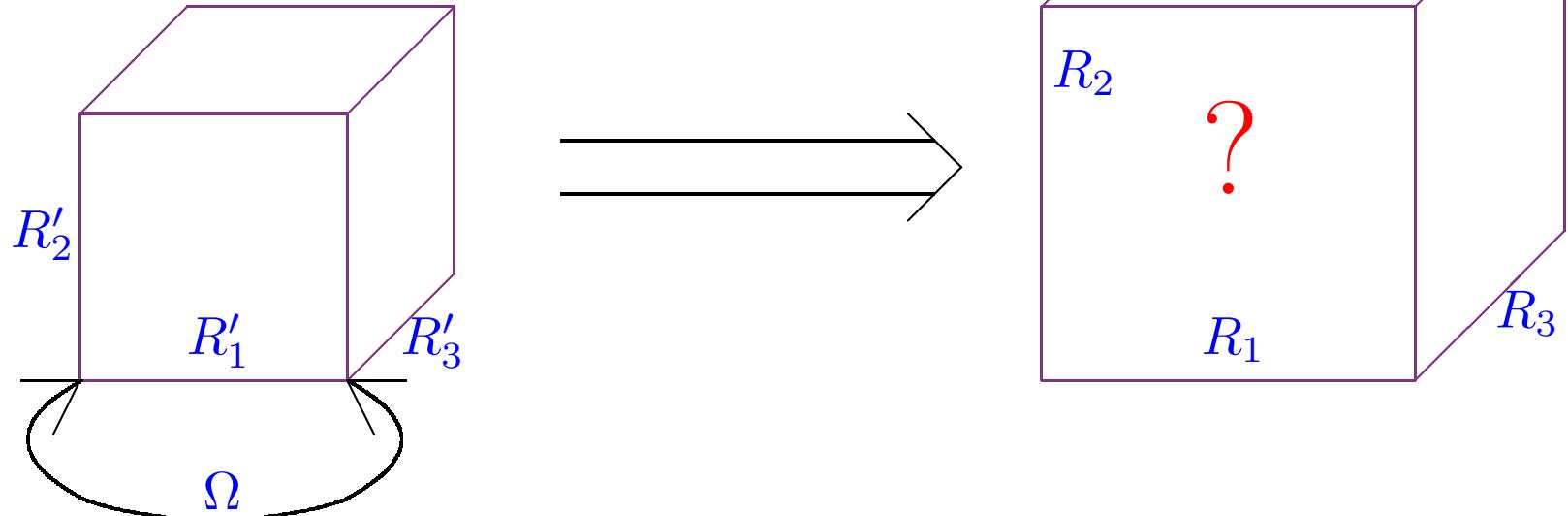
Causality: $v_{\max} = \sqrt{-\frac{g_{00}}{g_{xx}}} = (4\pi N)^{-1/2} g_s^{-3/2} \beta \rho^3.$

$$\Rightarrow E \geq \ell_P m v^{2/3} \beta^{-1/3}$$

D3-branes in Strong RR-flux



D3-branes in Strong RR-flux



$$f \equiv 1 + \frac{1}{g_s^2 \alpha'^4} r^2 \vec{n}^T \zeta^T \zeta \vec{n} \quad [\text{cf. Russo \& Tseytlin, 1995}]$$

$$ds^2 = f^{-\frac{1}{2}} \sum_{i=1}^3 dx_i^2 + f^{\frac{1}{2}} (dr^2 + r^2 d\Omega_5^2 - dt^2) - \frac{f^{-\frac{1}{2}} r^4}{g_s^2 \alpha'^4} (\vec{n}^T \zeta d\vec{n})^2,$$

$$C'_4 = \frac{1}{4g_s^2 \alpha'^2} f^{-1} r^2 \vec{n}^T \zeta d\vec{n} \wedge dx_1 \wedge dx_2 \wedge dx_3$$

Summary

- Lorentz violating $SO(3)$ -invariant QFT.
- Decoupled from Gravity.
- Nonlocal.
- Characteristic TJ coupling.

Open issues

- Fundamental formulation?
- Theories other than $\mathcal{N} = 4$ SYM?
- Tests?
- Relation to Little-String-Theory?
- cf. M-theory duals of (p, q) 5-branes.

[Witten, 1997]

More open issues

- Which objects have $c > 1$?
cf. nonrelativistic dispersion relations of
NCYM solitons [Landsteiner & Lopez &
Tytgat, Bak & Lee & Park, Hashimoto
& Itzhaki, 2000].

Even more open issues

- Relation to nonassociative structure and nongeometrical twists? [Ramgoolam, Medeiros & Ramgoolam, Bouwknegt & Hannabuss & Mathai, Shelton & Taylor & Wecht, Ellwood & Hashimoto, Nastase, . . . , 2004-6]
- Lightlike limit? cf. Brane probes of limits of Melvin twists. [. . . , Hashimoto & Sethi, Robins & Sethi, Hashimoto & Thomas, Sheikh-Jabbari, Alishahiha & Safarzadeh & Yavartanoo, Gimon & Hashimoto & Hubeny & Lunin & Rangamani, Alishahiha & OG, Varadarajan & OG]