Puff Field Theory

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What is String Theory?

Virginia Tech

Based on

- Aki Hashimoto, Sharon Jue, Bom Soo Kim, Anthony Ndirango and OG, "Aspects of Puff Field Theory," [arXiv:hep-th/0702030]
- OG, "A New Lorentz Violating Nonlocal Field Theory From String-Theory," [arXiv:hep-th/0609107]

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- c > 1?

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[Amelino-Camelia, Ellis, Mavromatos, Nanopoulos and Sarkar, 1998]

$$\frac{\Delta c}{c} < \left| \frac{E}{10^{17} \text{GeV}} \right|, \qquad 1 \text{MeV} < E < 17 \text{MeV}$$

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Preserve Lorentz invariant photon dispersion relation.







$$\begin{array}{c} & & & \\ & & & & \\ & & & \\ & & & & \\ & & & \\ & & & & \\ & & & \\ & & & &$$

Yang-Mills theory on a Noncommutative ℝ⁴ (NCYM)
 [Douglas & Hull, Connes & Douglas & Schwarz,...]

$$\begin{array}{c} & & \\ & & & \\ & & & \\ & & & \\ & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & &$$

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$$Q = \text{R-charge}$$

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Nonlocality - Volume

• In this talk:



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Deformations of N = 4 **SYM**

$$L = L_{N=4} + \zeta \mathcal{O}^{(\Delta+4)} + \cdots$$
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Theory	Deformation $\zeta \mathcal{O}$	Δ						
NCYM	$(1/2g_{YM}^2)\theta^{\mu\nu} \operatorname{Tr} \{F_{\mu\sigma}F_{\nu\tau}F^{\sigma\tau}\}$	2	$\zeta \rightarrow \theta$					
	$+F_{\sigma\tau}F^{\mu\nu}F_{\mu\nu}+\cdots\}$							
Dipole	$(1/2g_{YM}^2)L^{\mu}_{IJ}\operatorname{Tr}\{F_{\mu\nu}(\Phi^I D^{\nu}\Phi^J + \cdots)$		$\dot{c} \rightarrow L$					
	$+\cdots\}$	-	· -					
This talk	•••••	3						



Our Construction (1)



$$g_{s} = \alpha' R_{3}'^{-1} R_{2}'^{-1} \qquad R_{2} = \alpha' \frac{5}{4} g_{s}'^{\frac{1}{2}} R_{1}'^{-\frac{1}{2}} R_{2}'^{-1} R_{1} = \alpha' \frac{3}{4} g_{s}'^{-\frac{1}{2}} R_{1}'^{-\frac{1}{2}} \qquad R_{3} = \alpha' \frac{5}{4} g_{s}'^{\frac{1}{2}} R_{1}'^{-\frac{1}{2}} R_{3}'^{-1}$$

Our Construction (1) Type-IIB Type-IIA $R'_1 \longrightarrow 0$ $\alpha'^{-1/2}R'_2 \longrightarrow \text{finite}$ $\alpha'^{-1/2}R'_3 \longrightarrow \text{finite} R_2$ N KK's N D3's $T_{23} \circ S \circ T_1$ R'_2 U(N) SYM R_3 R'_3 R'_1 R_1 $\alpha'^{-1/2} R_k \longrightarrow \infty, \quad (k = 1, 2, 3)$ $g_s \longrightarrow \text{finite}$ $R_2/R_1, R_3/R_1 \longrightarrow \text{finite}$

Our Construction (2)



Scaling Ω

$$\begin{split} \Omega &= \exp\left\{\frac{2\pi}{R_1R_2R_3}\zeta\right\} \longrightarrow I, \qquad \zeta \to \text{finite} \\ & \left(\begin{array}{cc} \beta_1 & & \\ -\beta_1 & & \\ & \beta_2 & \\ & -\beta_2 & \\ & & -\beta_2 & \\ & & -\beta_3 & \end{array}\right) \in so(6) \\ & & SUSY: \quad \begin{cases} 0\} \quad \subset \quad su(2) \quad \subset \quad su(3) \quad \subset \quad so(6) \\ \mathcal{N} = 4 & \mathcal{N} = 2 & \mathcal{N} = 1 & \mathcal{N} = 0 \end{split}$$









Electric and Magnetic fluxes

$$V \equiv R_1 R_2 R_3,$$

$$\begin{split} \mathbf{P} &\equiv \sum_{i=1}^{3} \frac{k_{i}}{R_{i}} \hat{\mathbf{n}}_{i}, \qquad \mathbf{E} \equiv \sum_{i=1}^{3} \frac{e_{i}R_{i}}{2\pi V} \hat{\mathbf{n}}_{i}, \quad \mathbf{B} \equiv \sum_{i=1}^{3} \frac{m_{i}R_{i}}{2\pi V} \hat{\mathbf{n}}_{i}. \\ E &= 2j\beta \frac{M_{s}^{4}}{g_{s}} + \frac{2\pi^{2}V^{2}}{|NV+2j\beta|} \left(\frac{g_{YM}^{2}}{2\pi} \mathbf{E}^{2} + \frac{2\pi}{g_{YM}^{2}} \mathbf{B}^{2}\right) \\ &+ |\mathbf{P} - \frac{4\pi^{2}V^{2}}{|NV+2j\beta|} \mathbf{E} \times \mathbf{B}|. \\ \mathcal{N} &= 2 \Longrightarrow \beta \equiv \beta_{1} = \beta_{2}, \qquad \beta_{3} = 0. \end{split}$$

cf. NCSYM [Ho, Morariu & Zumino, Hofman & Verlinde, Konechny & Schwarz, Pioline & Schwarz,...1998-9].

Supergravity Dual

$$ds^{2} = \frac{R^{2}}{r^{2}}K^{-\frac{1}{2}}\left[dx^{2} + dy^{2} + dz^{2} - \left(dt - \frac{4\pi N}{r^{2}}\vec{n}^{T}\zeta d\vec{n}\right)^{2}\right] \\ + \frac{R^{2}}{r^{2}}K^{\frac{1}{2}}dr^{2} + R^{2}K^{\frac{1}{2}}d\Omega_{5}^{2} \\ C_{4}' = \frac{\pi N}{r^{4}}K^{-1}dt \wedge dx \wedge dy \wedge dz \\ - \frac{\pi N}{g_{s}\alpha'^{2}r^{6}}K^{-1}\vec{n}^{T}\zeta d\vec{n} \wedge dx \wedge dy \wedge dz, \\ K \equiv 1 + \frac{16\pi^{2}N^{2}}{r^{6}}\vec{n}^{T}\zeta^{T}\zeta\vec{n} \\ \vec{n} \in S^{5}, \qquad d\Omega_{5}^{2} = \sum_{I=1}^{6}dn_{I}^{2}, \qquad R^{4} \equiv 4\pi g_{s}N{\alpha'}^{2}$$

IR limit

$$ds^{2} = \frac{R^{2}}{r^{2}} [dr^{2} + dx^{2} + dy^{2} + dz^{2} - dt^{2}] + R^{2} d\Omega_{5}^{2}$$
$$+ \frac{8\pi N R^{2}}{r^{4}} \vec{n}^{T} \zeta d\vec{n} dt + \cdots$$
$$C_{4}' = \frac{\pi N}{r^{4}} dt \wedge dx \wedge dy \wedge dz$$
$$- \frac{\pi N}{g_{s} {\alpha'}^{2} r^{6}} \vec{n}^{T} \zeta d\vec{n} \wedge dx \wedge dy \wedge dz + \cdots$$

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$$L_{N=4} \longrightarrow L_{N=4} + \zeta \mathcal{O} + \cdots$$

Short multiplets

SYM Operator	desc	SUGRA	dim	spin	$SU(4)_R$
$\operatorname{Tr} X^k$	-	$h^lpha_lpha \;\; a_{lphaeta\gamma\delta}$	k	(0,0)	(0,k,0)
		• • •			
$\operatorname{Tr} F_+ \overline{\lambda} \lambda X^k$	$Q^3 \overline{Q}$	A_{\mulpha}	k+5	$\left(\frac{1}{2},\frac{1}{2}\right)$	(1,k,1)
•••		• • •			•••
$\operatorname{Tr} F_+ F^2 X^k$	$Q^4 \overline{Q}^2$	$A_{\mu u}$	k+6	(1, 0)	(0,k,0)
		• • •			•••
$\operatorname{Tr} F_+ F \overline{\lambda} \lambda X^k$	$Q^3 \overline{Q}^3$	$h_{\mulpha} ~~ a_{\mulphaeta\gamma}$	k+7	$\left(\frac{1}{2},\frac{1}{2}\right)$	(1,k,1)
•••		•••	•••	•••	•••
$\operatorname{Tr} F_+^2 F^2 X^k$	$Q^4 \overline{Q}^4$	$h^lpha_lpha = a_{lphaeta\gamma\delta}$	k+8	(0,0)	(0,k,0)

taken from [D'Hoker & Freedman, 2003]

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$\operatorname{Tr} F_+ F^2 X^k$	$Q^4 \overline{Q}^2$	$A_{\mu u}$	k+6	(1,0)	(0,k,0)
•••	• • •	•••		• • •	•••
$\operatorname{Tr} F_{+}F_{-}\overline{\lambda}\lambda$	$Q^3 \overline{Q}^3$	$h_{\mulpha} ~~ a_{\mulphaeta\gamma}$	7	$(\frac{1}{2},\frac{1}{2})$	(1, 0, 1)
•••	• • •	• • •	•••	• • •	• • •
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IR-limit

$$L = L_{N=4} + \zeta \mathcal{O}^{(\Delta+4)} + \cdots$$
$$\zeta \mathcal{O} \to \zeta^{\mu}_{AB} T_{\mu\nu} J^{\nu AB} + \zeta^{\mu}_{AB} \epsilon^{ABCDEF} \epsilon_{\mu\nu\sigma\tau} X^C \partial_{\nu} X^D \partial_{\sigma} X^E \partial_{\tau} X^F + \cdots$$
$$J^{\nu AB} \equiv \text{R-current}$$
$$T^{\mu\nu} \equiv \text{Energy-momentum tensor}$$

 $A, B, C, \dots = 1, \dots, 6.$







UV limit (1)



UV limit (1)



UV Limit (2)

T-duality on fiber: Type-IIB \Rightarrow M.

$$\ell_P^{-2} ds_M^2 = (4\pi N)^{-1/3} g_s^{-1} \rho^2 \left[\Delta^{-2/3} (dx^2 + dy^2 + dz^2) - \Delta^{1/3} dt^2 \right] + (4\pi N)^{2/3} \Delta^{1/3} \rho^{-2} d\rho^2 + (4\pi N)^{2/3} \Delta^{1/3} ds_B^2 + (4\pi N)^{-1/3} \Delta^{1/3} (g_s^{-1} d\xi^2 + g_s d\eta^2),$$

$$\ell_P^{-3}G_4 = 2\pi N\omega \wedge \omega + 2\omega \wedge d\eta \wedge d\xi + \frac{2}{3}(4\pi N)^{-2}g_s^{-3}\beta d\left(\frac{\rho^6}{\Delta}\right) \wedge dx \wedge dy \wedge dz,$$

$$\Delta \equiv 1 + (4\pi N)^{-1} g_s^{-3} \beta^2 \rho^6, \qquad (\rho \propto 1/r)$$



Extreme UV Limit (2)

$$ho
ightarrow \infty$$

- Curvature $\rightarrow 0$.
- Geodesically complete.
- Redshift in frequency $(g_{00} \rightarrow \infty)$
- Blueshift in wavelength $(g_{xx}, g_{yy}, g_{zz} \rightarrow 0)$

Degrees of Freedom

UV regime of PFT – holographically dual to weakly-coupled background (?)

Is the spectrum discrete?

Degrees of Freedom

UV regime of PFT – holographically dual to weakly-coupled background (?)

Is the spectrum discrete?

For particle trajectories with fixed (E, \vec{p}) :

$$\rho \le \rho_{\max} \equiv (4\pi N)^{1/6} g_s^{1/2} \beta^{-1/3} \left(\frac{E}{|\vec{p}|}\right)^{1/3}$$

Suggests a discrete spectrum!

cf. Little-String-Theory. [Aharony & Berkooz & Kutasov & Seiberg, 1998]

$$L = L_{(N=4)} + \zeta(T^{00}J_0 + \dots) + \dots$$

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• In R-charged matter:

$$g_{00} \to g_{00} - \zeta \langle J^0 \rangle + \cdots, \qquad v_{\max} \approx 1 + \frac{1}{2} \zeta \langle J^0 \rangle.$$

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• In dense matter:

$$\Delta V_{\text{eff}} \approx \zeta \langle T^{00} \rangle Q^{(R)}, \qquad Q^{(R)} \equiv \int J^0 d^3 x = (\text{R-charge}).$$

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Trap R-charged particles?

More on c > 1

Extreme UV:

$$\ell_P^{-2} ds_M^2 \approx -(4\pi N)^{-2/3} g_s^{-2} \beta^{2/3} \rho^4 dt^2 +(4\pi N)^{1/3} g_s \beta^{-4/3} \rho^{-2} (dx^2 + dy^2 + dz^2) +(4\pi N)^{1/3} g_s^{-1} \beta^{2/3} d\rho^2 + \cdots$$

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For a particle of mass m:

$$-\ell_P^2 m^2 \ge -(4\pi N)^{2/3} g_s^2 \beta^{-2/3} \frac{E^2}{\rho^4}.$$

Causality:
$$v_{\max} = \sqrt{-\frac{g_{00}}{g_{xx}}} = (4\pi N)^{-1/2} g_s^{-3/2} \beta \rho^3.$$

 $E \ge \ell_P m v^{2/3} \beta^{-1/3}$





Summary

- Lorentz violating SO(3)-invariant QFT.
- Decoupled from Gravity.
- Nonlocal.
- Characteristic TJ coupling.

Open issues

- Fundamental formulation?
- Theories other than $\mathcal{N} = 4$ SYM?
- Tests?
- Relation to Little-String-Theory?
- cf. M-theory duals of (p, q) 5-branes. [Witten, 1997]

More open issues

• Which objects have c > 1?

cf. nonrelativistic dispersion relations of NCYM solitons [Landsteiner & Lopez & Tytgat, Bak & Lee & Park, Hashimoto & Itzhaki, 2000].

Even more open issues

- Relation to nonassociative structure and nongeometrical twists? [Ramgoolam, Medeiros & Ramgoolam, Bouwknegt & Hannabuss & Mathai, Shelton & Taylor & Wecht, Ellwood & Hashimoto, Nastase, ..., 2004-6]
- Lightlike limit? cf. Brane probes of limits of Melvin twists. [..., Hashimoto & Sethi, Robins & Sethi, Hashimoto & Thomas, Sheikh-Jabbari, Alishahiha & Safarzadeh & Yavartanoo, Gimon & Hashimoto & Hubeny & Lunin & Rangamani, Alishahiha & OG, Varadarajan & OG]